

THE HUBBARD MODEL AND BEYOND

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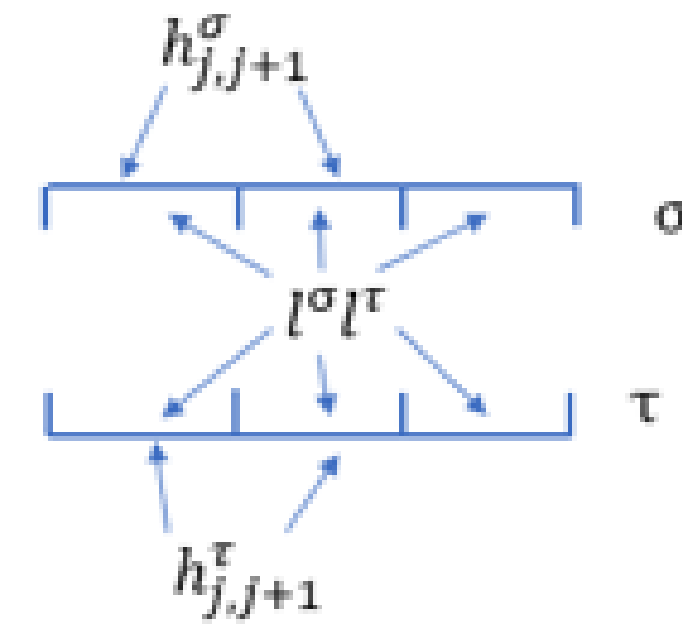
First range 3 elliptic deformation of the Hubbard model

$$\mathbb{H}_3 = \sum_j h_{j,j+1}^\sigma + h_{j,j+1}^\tau + U l_{j,j+1,j+2}^\sigma l_{j,j+1,j+2}^\tau,$$

$$h_{j,j+1} = i(\sigma_j^+ \sigma_{j+1}^- - \sigma_j^- \sigma_{j+1}^+)$$

$$l_{j,j+1,j+2} = Z_{j+1} + \kappa (X_j + X_{j+2}) X_{j+1} - \kappa^2 X_j Z_{j+1} X_{j+2}$$

U coupling constant κ deformation parameter



Where is the Hubbard model?

$\kappa = 0$ Hubbard model

Jordan-Wigner transformation $\sigma \rightarrow c^\dagger$ $\tau \rightarrow c^\dagger$

$$\mathbb{H}_{Hub} = \sum_j \left[(c_j^\dagger)^\dagger c_{j+1}^\dagger + (c_{j+1}^\dagger)^\dagger c_j^\dagger + (c_j^\dagger)^\dagger c_{j+1}^\dagger + (c_{j+1}^\dagger)^\dagger c_j^\dagger + \frac{U}{4} (1 - 2n_j^\uparrow)(1 - 2n_j^\downarrow) \right],$$

$n = c^\dagger c$, c^\dagger , c fermionic creation and annihilation operators

Corresponding range 2 model

Bond-site transformation $X_j \rightarrow X_{j-\frac{1}{2}} X_{j+\frac{1}{2}}$, $Z_i Z_{i+1} \rightarrow Z_{i+1/2}$

$$\begin{cases} \uparrow\uparrow \\ \downarrow\downarrow \end{cases} \rightarrow \uparrow \quad \begin{cases} \uparrow\downarrow \\ \downarrow\uparrow \end{cases} \rightarrow \downarrow$$

$\kappa = \tan \frac{\theta}{2}$

$$\mathbb{H}_2 = \sum_j [h_{j,j+1}^\sigma + h_{j,j+1}^\tau + u L_{j,j+1}^\sigma L_{j,j+1}^\tau]$$

$$h_{j,j+1} = \frac{dn}{k} (\sigma_j^+ \sigma_{j+1}^- - \sigma_j^- \sigma_{j+1}^+), \quad L_{j,j+1} = \begin{pmatrix} \sin \theta & 0 & 0 & \cos \theta \\ 0 & 0 & 1 & 0 \\ 0 & 1 & 0 & 0 \\ \cos \theta & 0 & 0 & -\sin \theta \end{pmatrix}$$

Open quantum system interpretations: general mapping

Map between Hermitian \mathbb{H} and Lindblad superoperator

$$\mathbb{H} = \sum_j h^\sigma + h^\tau + U l^\sigma l^\tau$$

$$\text{Conditions: } h = -h^T, \quad l = l^* = l^T, \quad l^2 = l^\dagger l = \mathbb{I}$$

Complexify $U \rightarrow iU$

$$\mathbb{H} \rightarrow \mathcal{L} = -ih_{j,j+1}^\sigma + ih^{\tau T} + U \left(l^\sigma l^{\tau*} - \frac{1}{2} l^{\sigma\dagger} l^\sigma - \frac{1}{2} l^{\tau T} l^{\tau*} \right),$$

Integrability

Range 2 model

R -matrix : unusual functional dependence $\sin \frac{1}{2} [\text{am}(u|k^2) - \text{am}(v|k^2)]$

Range 3 model

$$\check{R}\text{-matrix (2 site)} \xrightarrow{\text{bond-site}} \text{Lax operator (3 site)} \xrightarrow{\text{RLL}} \check{R}\text{-matrix (4 site)}$$

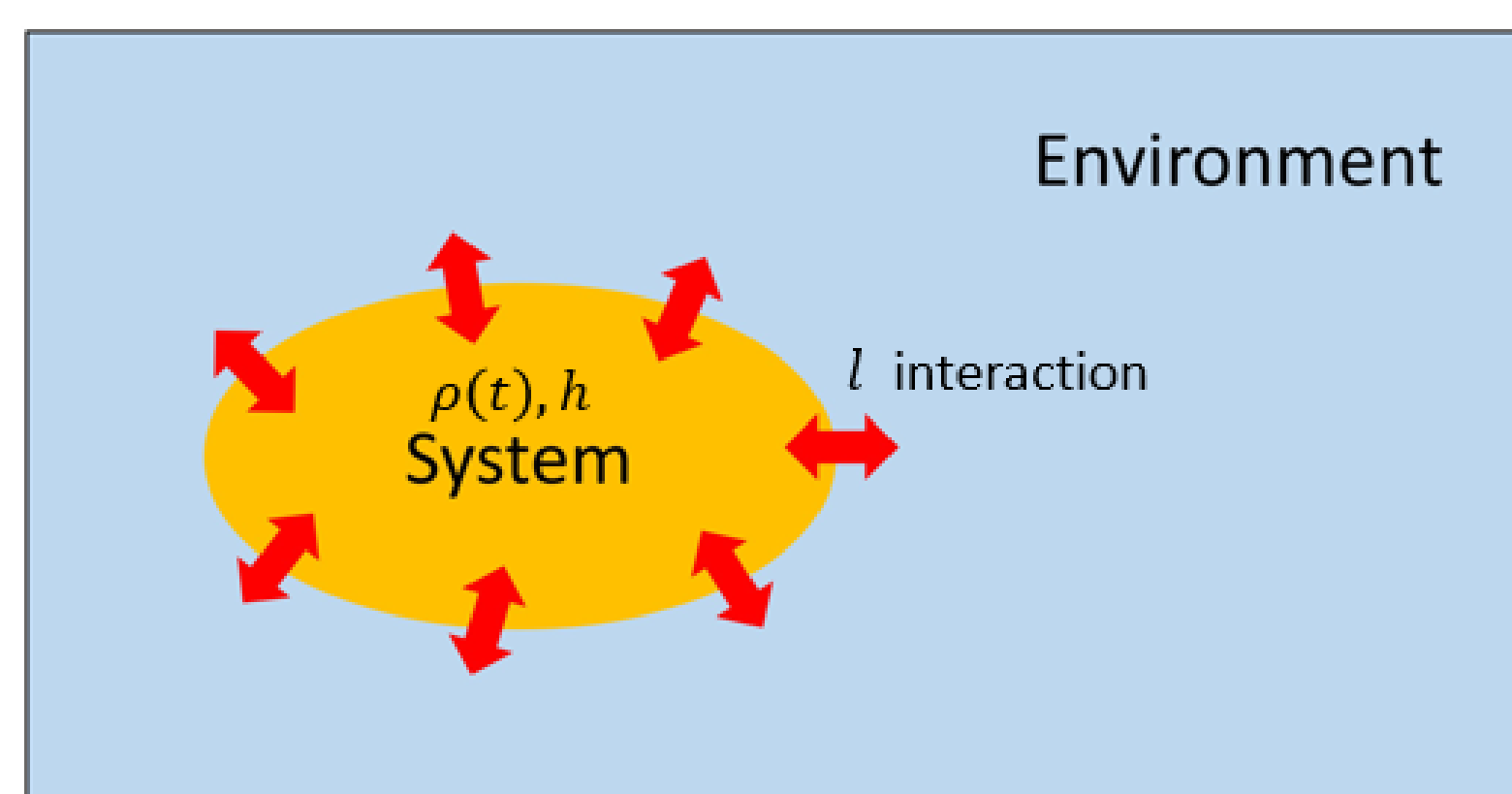
$$\check{R}_{34,56}(u_1, u_2) \check{R}_{12,34}(u_1, u_3) \check{R}_{34,56}(u_2, u_3) = \check{R}_{12,34}(u_2, u_3) \check{R}_{34,56}(u_1, u_3) \check{R}_{12,34}(u_1, u_2)$$

Lindblad superoperator

Approximation:

- Markovian evolution

$$\tau_{\text{environment}} \ll \tau_{\text{system}}$$

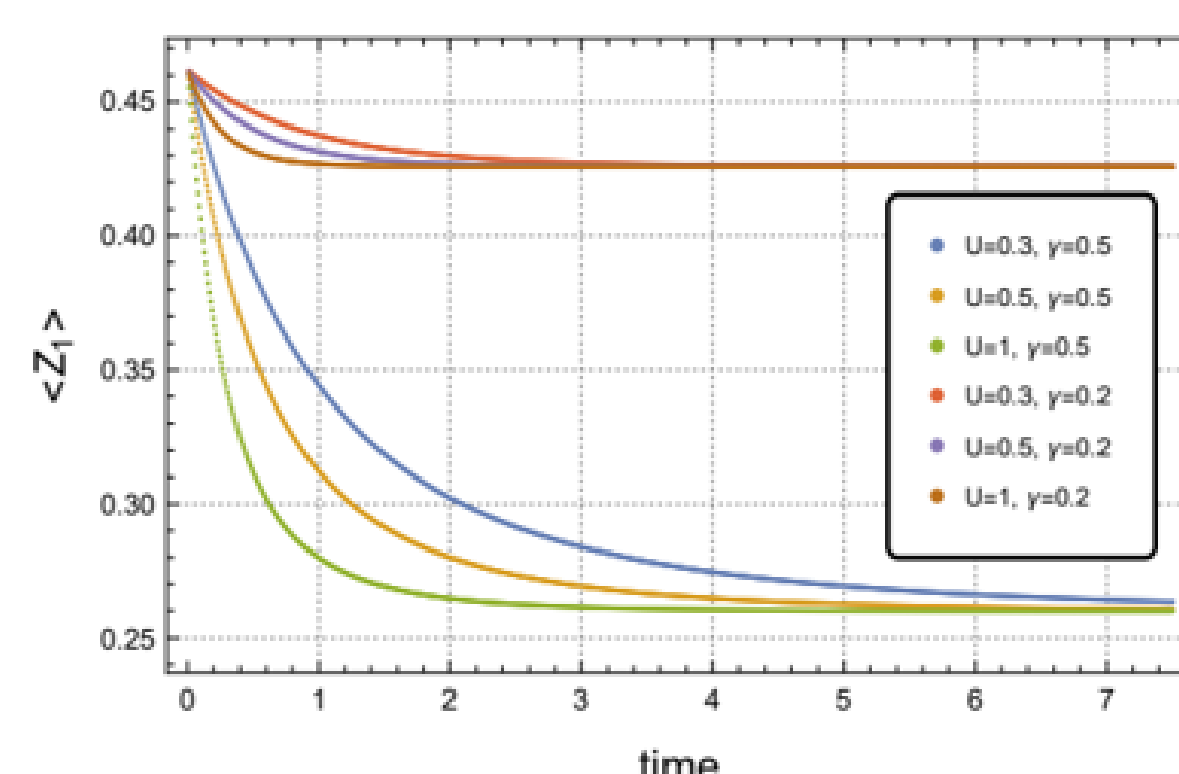


$$\dot{\rho}(t) = \mathcal{L}\rho(t) \rightarrow \underbrace{\dot{\rho}(t) = i[\rho, h]}_{\text{von-Neumann equation}} + U \underbrace{\left[\ell \rho \ell^\dagger - \frac{1}{2} \{ \ell^\dagger \ell, \rho \} \right]}_{\text{dissipator}}$$

Physical properties

Hidden "strong symmetries" Q : $[Q, h] = [Q, l] = 0$

	Q	# NESS	NESS
$\kappa = 0$	$Q_0 = \sum_i Z_i$	$L + 1$	$\rho_0 \sim e^{\alpha Q_0}$
$\kappa \neq 0$	$Q_\kappa = T(\kappa)^\dagger Q_0 T(\kappa)$	$L + 1$	$\rho_\kappa \sim T(\kappa)^\dagger \rho_0 T(\kappa)$



$$T(\kappa) = \text{Tr}_A(A_L(\kappa) A_{L-1}(\kappa) \dots A_1(\kappa)), \quad [T(\kappa), T(\kappa')] = 0$$

Non-unique NESS: memory - Long time limit of observable depend on initial state

Remarks

Why is this model new?

- First integrable deformation of Hubbard spanning 3 sites
- Unusual functional dependence in the R -matrix
- κ terms break $U(1)$ symmetry: XYZ-type deformation

Another deformation

Model B2 of PRL 126, 240403 (2021)

$$\mathbb{H}_{B2} = \sum_j h_{j,j+1}^\sigma + h_{j,j+1}^\tau + U l_{j,j+1}^\sigma l_{j,j+1}^\tau,$$

$$h_{j,j+1} = i(\sigma_j^+ \sigma_{j+1}^- - \sigma_j^- \sigma_{j+1}^+)$$

$$l_{j,j+1} = Z_j + \eta (X_j X_{j+1} + Y_j Y_{j+1}) + \eta^2 Z_{j+1}$$

$\eta = 0$ Hubbard model

(1)

[Murakami '98, Medvedyeva Essler Prosen '16, Ziolkowska Essler '19]

Open questions

- Solution of the model: Nested Bethe Ansatz + XYZ-methods
- Construction of integrable spin chains with open boundary conditions
- Implication of the Hermitian-Lindbladian mapping in both integrable and chaotic case
- Analytical computation of Liouville gap