



# Radiation and Reaction at One Loop

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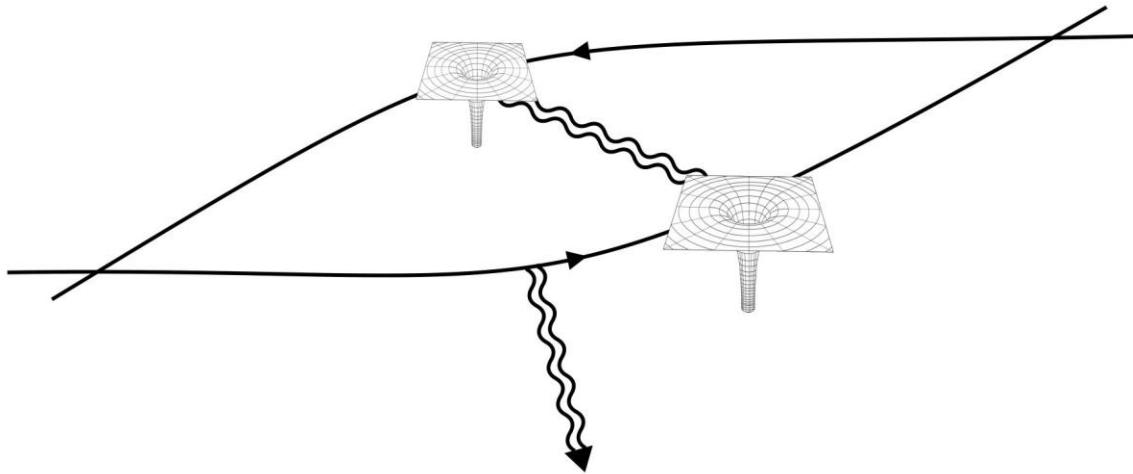
Higgs Centre for Theoretical Physics, The University of Edinburgh.

Based on “Radiation and Reaction at One Loop”  
With Donal O’Connell & Ingrid A. Vazquez-Holm.

arXiv: SOON.

QCD meets Gravity 2022, 16/12/2022

# Background and motivation



$$\Rightarrow \langle R_{\mu\nu\rho\sigma}(x) \rangle = \int_k \mathcal{A}(k) e^{-ik \cdot x}$$

- Kosower, Maybee & O'Connell, 2018  
Kosower, O'Connell, Gonzo & Cristofoli, 2019  
Sergola, Peinador-Veiga, Monteiro & O'Connell, 2021  
Sergola, White, Moynihan, Ross, Cristofoli, Gonzo & O'Connell, 2021

See also  
Graham's and  
Radu's talks!

# Classical radiation at one loop

$$\langle F_{\mu\nu}(x) \rangle = \sum_{\eta} \int d\Phi(k) \left( k_{[\mu} \varepsilon_{\nu]}^{\eta} \langle \psi | S^\dagger a_{\eta}(k) S | \psi \rangle e^{-ik \cdot x} + \text{h.c.} \right)$$

$|\psi\rangle \sim \int |p_1, p_2\rangle$   
↓      ↓  
↑      {  
 On-shell       $\alpha_{\eta}(k)$   
 integral       $S = 1 + iT$   
 support

Shen, 2018

Kosower, Maybee & O'Connell, 2018

Kosower, O'Connell, Gonzo & Cristofoli, 2019

Vazquez-Holm & Carrasco, 2020

Vazquez-Holm & Carrasco, 2021

Herrmann, Parra-Martinez, Ruf & Zeng, 2021

# Classical radiation at one loop

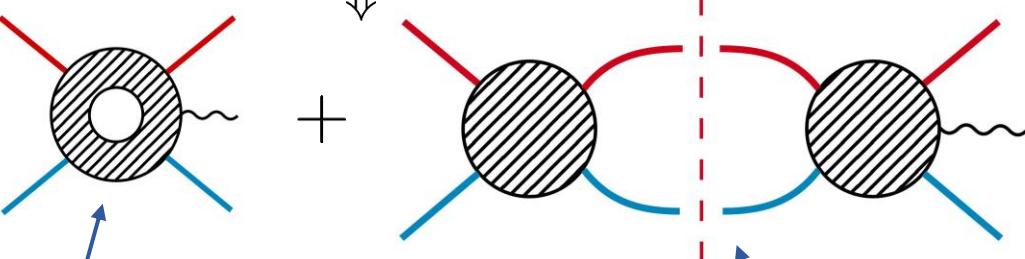
$$\langle F_{\mu\nu}(x) \rangle = \sum_{\eta} \int d\Phi(k) \left( k_{[\mu} \varepsilon_{\nu]}^{\eta} \langle \psi | S^\dagger a_{\eta}(k) S | \psi \rangle e^{-ik \cdot x} + \text{h.c.} \right)$$

$|\psi\rangle \sim \int |p_1, p_2\rangle$   
 $\alpha_{\eta}(k)$   
 $S = 1 + iT$

On-shell integral support

$\alpha_{\eta}(k) = i \mathcal{A}_{1L} + \mathcal{A}_{tree} \times \mathcal{A}_{tree}^*$

Taken in the limit  $\hbar \rightarrow 0$



Shen, 2018

Kosower, Maybee & O'Connell, 2018

Kosower, O'Connell, Gonzo & Cristofoli, 2019

Vazquez-Holm & Carrasco, 2020

Vazquez-Holm & Carrasco, 2021

Herrmann, Parra-Martinez, Ruf & Zeng, 2021



# Classical radiation at one loop

$$\alpha_\eta(k) = \frac{1}{2} \langle \psi | i a_\eta(k) (T + T^\dagger) + T^\dagger [a_\eta(k), T] - [a_\eta(k), T^\dagger] T | \psi \rangle$$

Real part

Imaginary part

A diagram illustrating the decomposition of the classical radiation term. The term is enclosed in a large blue bracket at the bottom. Above it, a smaller blue bracket groups the first two terms:  $i a_\eta(k) (T + T^\dagger)$ . This is labeled "Real part". Below the large bracket, another blue bracket groups the last two terms:  $T^\dagger [a_\eta(k), T] - [a_\eta(k), T^\dagger] T$ . This is labeled "Imaginary part".

Herrmann, Parra-Martinez, Ruf & Zeng, 2021  
Bern, Parra-Martinez, Roiban, Ruf, Shen, Solon  
& Zeng 2021

# Classical radiation at one loop

$$\text{Re} \quad \text{Diagram} = \text{Diagram}_1 + \text{Diagram}_2 \sim g^5 \int d^4\ell N i(P - P^*)$$

↑  
Real part       $\frac{1}{p^2 - m^2 + i\epsilon} = \text{P.V.} \frac{1}{p^2 - m^2} + i\delta(p^2 - m^2)$

$$\alpha_\eta(k) = \frac{1}{2} \langle \psi | i a_\eta(k) (T + T^\dagger) + T^\dagger [a_\eta(k), T] - [a_\eta(k), T^\dagger] T | \psi \rangle$$

Imaginary part

# Classical radiation at one loop

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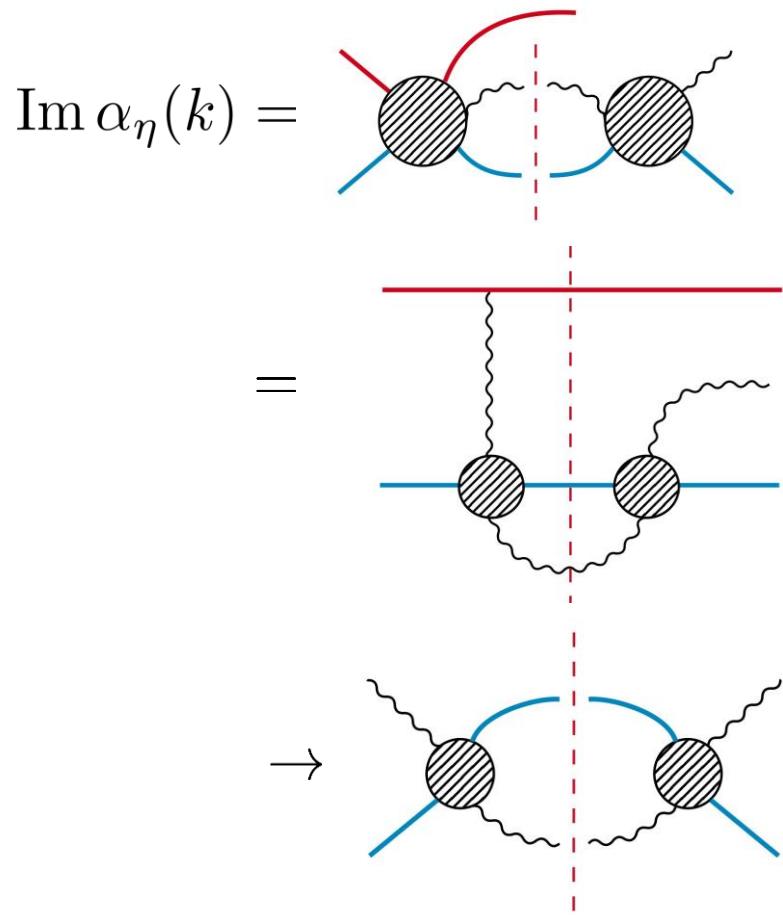
$$\alpha_\eta(k) = \frac{1}{2} \langle \psi | i a_\eta(k) (T + T^\dagger) + T^\dagger [a_\eta(k), T] - [a_\eta(k), T^\dagger] T | \psi \rangle$$

Imaginary part

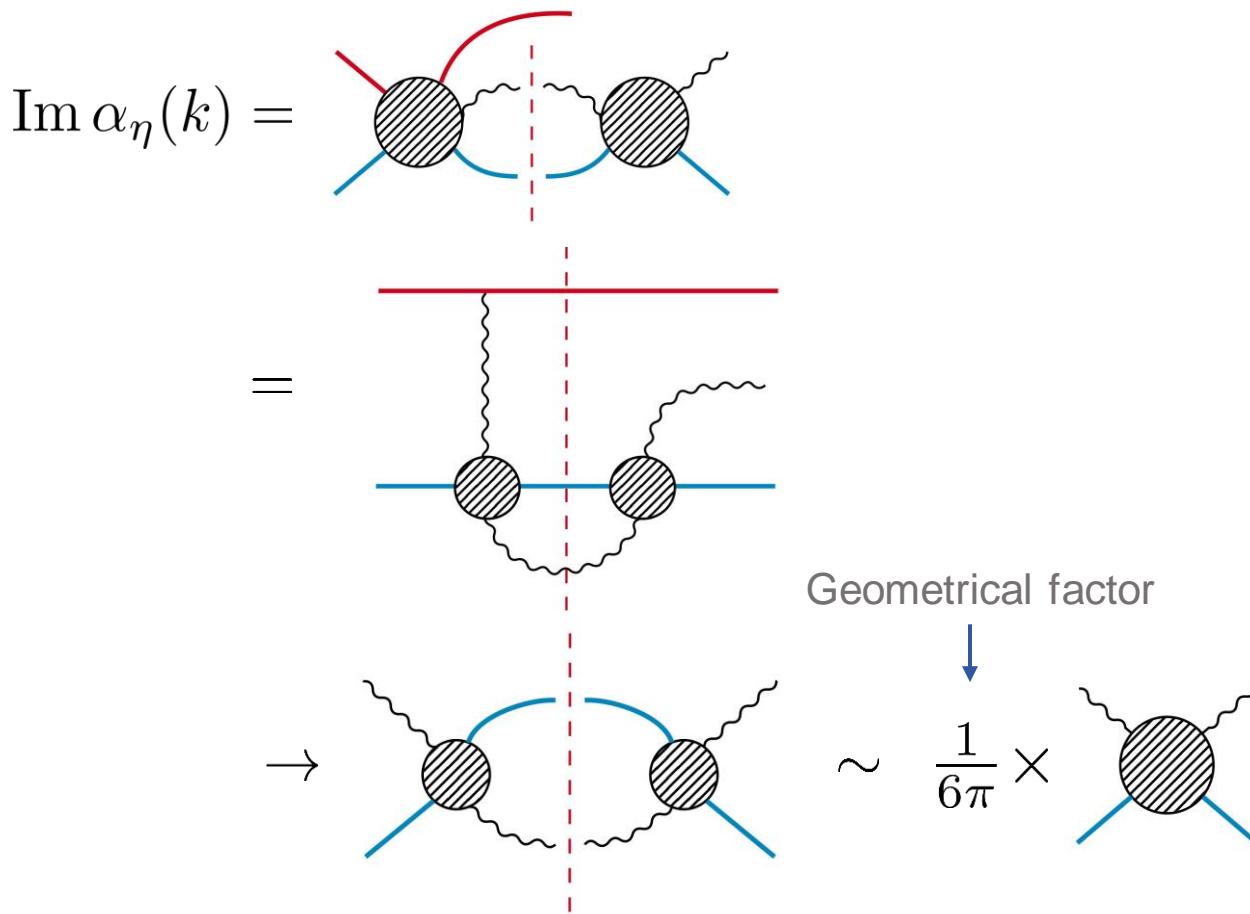
$$\text{Diagram}_1 + \text{Diagram}_2 + \text{exchanges}$$

Herrmann, Parra-Martinez, Ruf & Zeng, 2021  
 Bern, Parra-Martinez, Roiban, Ruf, Shen, Solon  
 & Zeng 2021

# Radiation reaction and imaginary part



# Radiation reaction and imaginary part





# QED example: ALD force

$$m_2 \frac{dp_2^\mu}{d\tau} = Q F^{\mu\nu} p_\nu + \frac{Q^2}{6\pi} \left( \frac{d^2 p_2^\mu}{d\tau^2} + \frac{p_2^\mu}{m_2^2} \frac{dp_2}{d\tau} \cdot \frac{dp_2}{d\tau} \right)$$

Lorentz force

Radiation reaction force

“Schott term”

$$\Rightarrow x^\mu(\tau) = b^\mu + u^\mu \tau + \dots + Q^4 \delta x_{\text{ALD}}^\mu(\tau)$$

# QED example: ALD force

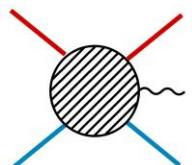
$$m_2 \frac{dp_2^\mu}{d\tau} = Q F^{\mu\nu} p_\nu + \frac{Q^2}{6\pi} \left( \frac{d^2 p_2^\mu}{d\tau^2} + \frac{p_2^\mu}{m_2^2} \frac{dp_2}{d\tau} \cdot \frac{dp_2}{d\tau} \right)$$

Lorentz force

Radiation reaction force

“Schott term”

$$\Rightarrow x^\mu(\tau) = b^\mu + u^\mu \tau + \dots + Q^4 \delta x_{\text{ALD}}^\mu(\tau)$$

$$\Rightarrow -k^2 \tilde{A}^\mu(k) = \dots + \frac{Q^5}{6\pi} \times$$


ALD, 1892-1905. Schott, 1912.  
 Kosower, Maybee & O'Connell, 2018

# Radiation in QCD

$$\langle F_{\mu\nu}^A(x) \rangle = \sum_{\eta} \int d\Phi(k) \left( -ik_{[\mu} \varepsilon_{\nu]}^{\eta} \langle \psi | S^\dagger a_{\eta}^A(k) S | \psi \rangle e^{-ik \cdot x} + \text{h.c.} \right)$$

“QED-like” RR

$$\begin{aligned}
 \mathcal{A}_{1L} = & \mathcal{C} \left( \begin{array}{c} \text{red} \\ \diagup \quad \diagdown \\ \text{blue} \end{array} \right) A_1 + \mathcal{C} \left( \begin{array}{c} \text{red} \\ \diagdown \quad \diagup \\ \text{blue} \end{array} \right) A_2 + \mathcal{C} \left( \begin{array}{c} \text{red} \\ \diagup \quad \diagup \\ \text{blue} \end{array} \right) A_3 + \mathcal{C} \left( \begin{array}{c} \text{red} \\ \diagup \quad \diagup \\ \text{blue} \end{array} \right) A_4 \\
 & + \mathcal{C} \left( \begin{array}{c} \text{red} \\ \diagup \quad \diagup \\ \text{blue} \end{array} \right) A_5 + \mathcal{C} \left( \begin{array}{c} \text{red} \\ \diagup \quad \diagup \\ \text{blue} \end{array} \right) A_6 + \mathcal{C} \left( \begin{array}{c} \text{red} \\ \diagup \quad \diagup \\ \text{blue} \end{array} \right) A_7 + \mathcal{C} \left( \begin{array}{c} \text{red} \\ \diagup \quad \diagup \\ \text{blue} \end{array} \right) A_8 \\
 & + \dots
 \end{aligned}$$

Non-Abelian RR?

# Maximally non-Abelian sector: $A_6$



$$A_6 = \begin{array}{c} p_1 \xrightarrow[q_1]{\text{---}} p_1 + q_1 \\ \text{---} \quad \text{---} \\ \text{---} \quad \text{---} \end{array} + \dots$$

+  $\cdots$

$$= \frac{1}{q_1^2} \int d^4 \ell \delta(p_2 \cdot \ell) \varepsilon \cdot \sum_i J_i$$



$$J_1^\mu = -\frac{\ell^\mu}{(\ell + q_2)^2 (\ell - q_1)^2} \left( \frac{p_1 \cdot p_2}{4 p_2 \cdot q_1} + \frac{\ell \cdot p_1}{\ell^2} \right) p_2^2,$$

$$J_2^\mu = \frac{q_1^\mu}{\ell^2 (\ell + q_2)^2 (\ell - q_1)^2} (p_2^2 \ell \cdot p_1 + p_1 \cdot p_2 p_2 \cdot q_1),$$

$$J_3^\mu = \frac{p_1^\mu}{\ell^2} \left( \frac{p_2^2}{4(\ell + q_2)^2} - \frac{(p_2 \cdot q_1)^2}{(\ell + q_2)^2 (\ell - q_1)^2} \right),$$

$$J_4^\mu = \frac{p_2^\mu}{\ell^2} \left( -\frac{p_2^2 \ell \cdot p_1}{4(\ell - q_1)^2 p_2 \cdot q_1} - \frac{p_2^2 p_1 \cdot p_2}{4(p_2 \cdot q_1)^2} + \frac{(p_1 \cdot p_2 q_1 \cdot q_2 - q_1 \cdot p_2 q_2 \cdot p_1)}{(\ell + q_2)^2 (\ell - q_1)^2} \right)$$

# Maximally non-Abelian sector: $A_6$

$$\text{Im } A_6 = \begin{array}{c} p_1 \text{ --- } p_1 + q_1 \\ | \qquad \qquad | \\ q_1 \text{ --- } \text{ (wavy line) } \text{ (wavy line) } \\ | \qquad \qquad | \\ p_2 \text{ --- } p_2 + q_2 \end{array} + \dots$$

$$= \frac{1}{q_1^2} \int d^4\ell \delta(p_2 \cdot \ell) \delta((q_1 - \ell)^2) \varepsilon \cdot \sum_i J_i$$



$$J_1^\mu = -\frac{\ell^\mu}{(\ell + q_2)^2 (\ell - q_1)^2} \left( \frac{p_1 \cdot p_2}{4 p_2 \cdot q_1} + \frac{\ell \cdot p_1}{\ell^2} \right) p_2^2,$$

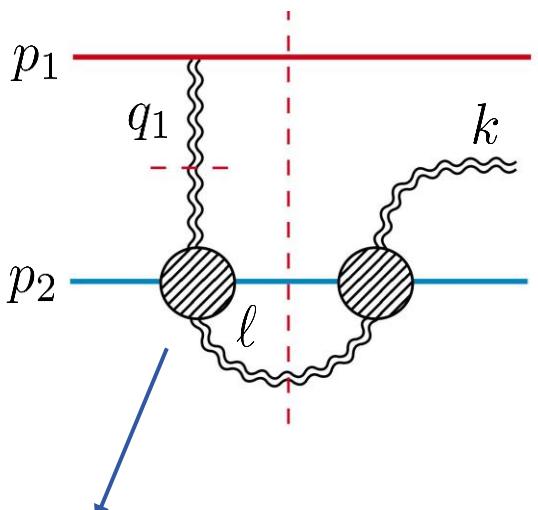
$$J_2^\mu = \frac{q_1^\mu}{\ell^2 (\ell + q_2)^2 (\ell - q_1)^2} (p_2^2 \ell \cdot p_1 + p_1 \cdot p_2 p_2 \cdot q_1),$$

$$J_3^\mu = \frac{p_1^\mu}{\ell^2} \left( \frac{p_2^2}{4(\ell + q_2)^2} - \frac{(p_2 \cdot q_1)^2}{(\ell + q_2)^2 (\ell - q_1)^2} \right),$$

$$J_4^\mu = \frac{p_2^\mu}{\ell^2} \left( -\frac{p_2^2 \ell \cdot p_1}{4(\ell - q_1)^2 p_2 \cdot q_1} - \frac{p_2^2 p_1 \cdot p_2}{4(p_2 \cdot q_1)^2} + \frac{(p_1 \cdot p_2 q_1 \cdot q_2 - q_1 \cdot p_2 q_2 \cdot p_1)}{(\ell + q_2)^2 (\ell - q_1)^2} \right)$$

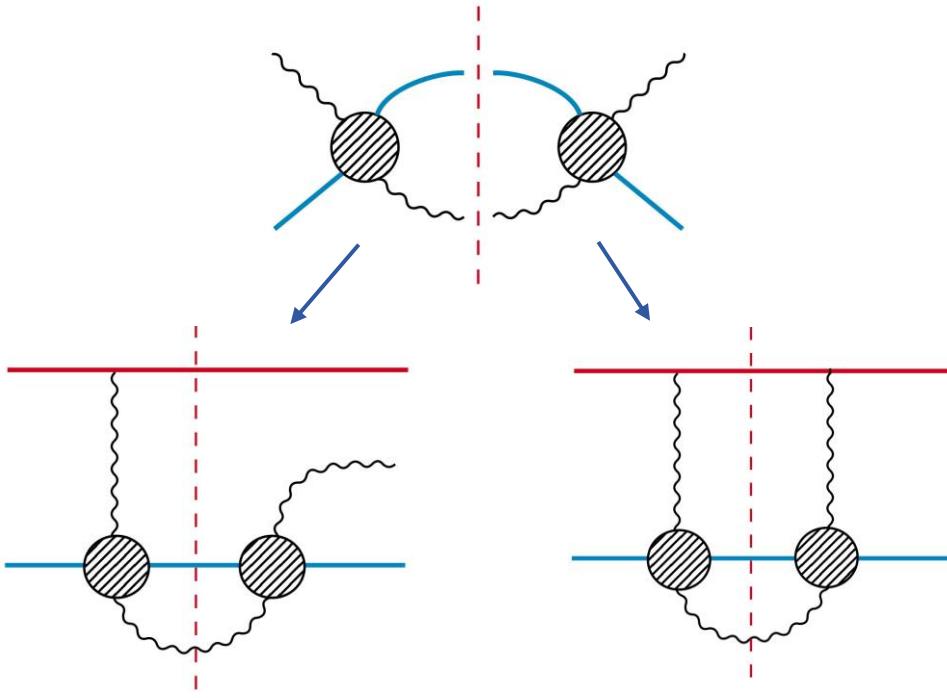
# Gravity imaginary part

$$\text{Im } \mathcal{A}_{\text{GR}} = \frac{1}{q_1^2} \int d\Phi(l) \delta(p_2 \cdot (q_1 - l)) \frac{(p_1 \cdot k)^4}{q_1 \cdot l \ k \cdot l} \times \sum_{\text{helicities}} (\varepsilon \cdot \varepsilon)^2 (\varepsilon \cdot \varepsilon)^2$$



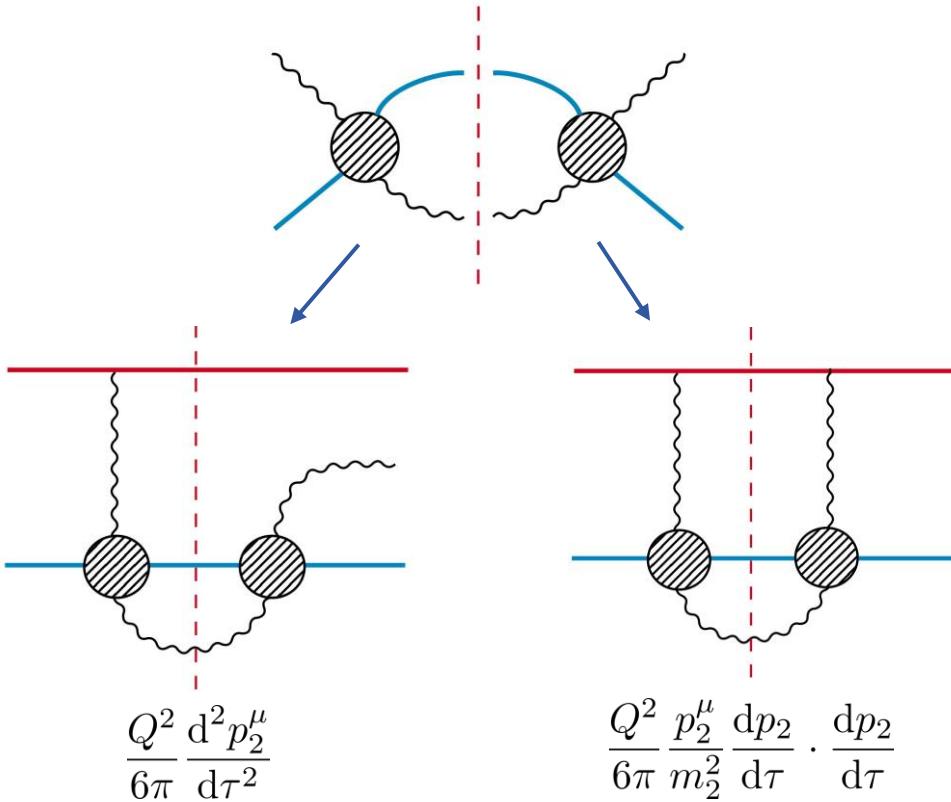
$$\mathcal{A}_{\text{GR}}^{\text{Compt.}} \propto \left( \mathcal{A}_{\text{EM}}^{\text{Compt.}} \right)^2$$

# Cut universality



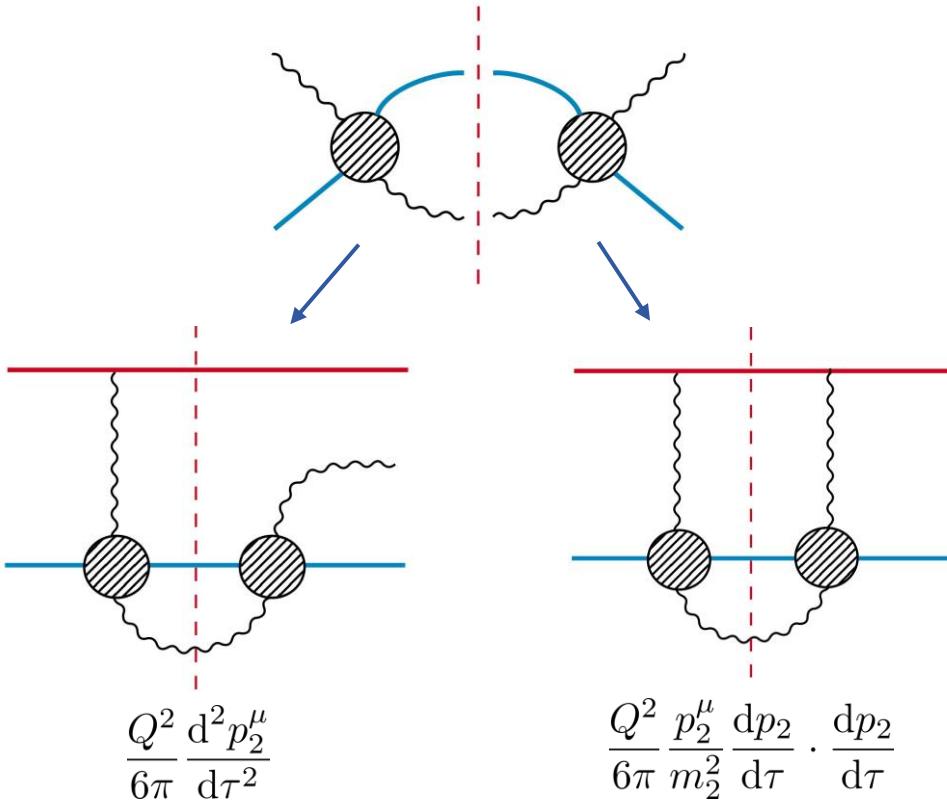
Mino, Sasaki & Tanaka, 1997  
Quinn & Wald, 1997  
Kosower, Maybee & O'Connell, 2018  
Dlapa, Kalin, Liu, Neef & Porto, 2022

# Cut universality



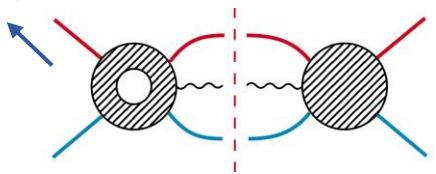
Mino, Sasaki & Tanaka, 1997  
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# Cut universality

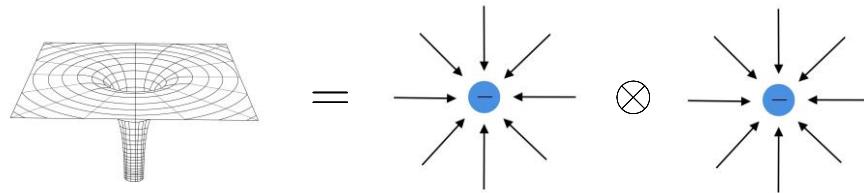


Mino, Sasaki & Tanaka, 1997  
 Quinn & Wald, 1997  
 Kosower, Maybee & O'Connell, 2018  
 Dlapa, Kalin, Liu, Neef & Porto, 2022

$$\Delta p_{3L}^\mu = \int q^\mu \mathcal{A} + \int \ell^\mu \mathcal{A}^* \mathcal{A}$$



# Gravity



$$R_{\mu\nu\rho\sigma}(x) = -\sum_{\eta} \int d\Phi(k) \left( k_{[\mu} \varepsilon_{\nu]}^{\eta} k_{[\rho} \varepsilon_{\sigma]}^{\eta} \alpha_{\eta}(k) e^{-ik \cdot x} + \text{h.c.} \right)$$

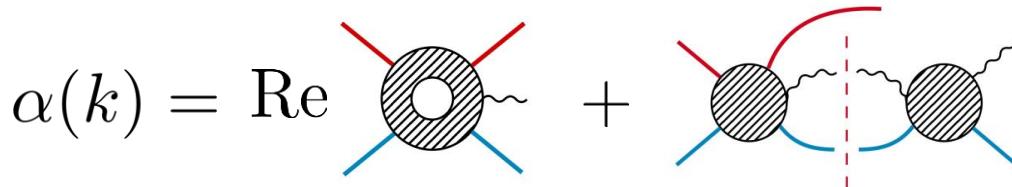
↑

Compute using the DC of  
Vazquez-Holm & Carrasco, 2021

- Bern, Carrasco & Johansson, 2008  
Goldberger & Ridgway, 2016  
Luna, Nicholson O'Connell & White, 2017  
Vazquez-Holm & Carrasco, 2020  
Vazquez-Holm & Carrasco, 2021

# Conclusions and future directions

- Amplitudes for classical radiation:

$$\alpha(k) = \text{Re} \quad \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \end{array} + \quad \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \end{array}$$


The diagram illustrates the definition of the amplitude  $\alpha(k)$ . It consists of two parts separated by a plus sign. The first part is the real part of a one-loop Feynman diagram, which is a shaded circle with a wavy line entering from the left and a blue line exiting to the right. The second part is a two-loop diagram where two shaded circles are connected by a red curved line and a blue wavy line, with a vertical dashed red line passing through the center of the circles.

- We care about gravitational waveforms, it is natural use BCJ numerators.
- The physics is surprisingly rich.
- Things seem to be universal from an amplitude perspective.
- *Next:* Compute gravity waveform.



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**Thank you!**