

All-orders celestial OPEs from the worldsheet

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14 Dec, 2022

QCD meets gravity 2022

Based on arxiv:2111.02279, 2111.15584, 2211.17124
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Quick background on celestial holography

Codimension 2 correspondence between quantum gravity in 4d Minkowski space and some conformal field theory on 2d celestial sphere.

$$\text{Observation: } \mathfrak{sl}(2, \mathbb{C}) \cong \mathfrak{so}(3, 1)$$

Conformal primary operators $\mathcal{O}_{h, \bar{h}}(z, \bar{z})$ in usual 2d CFTs diagonalizes dilation ($\Delta = h + \bar{h}$) and rotation ($s = h - \bar{h}$).

The analogous object in 4d diagonalizes Lorentz boost and Lorentz rotation. For massless particles, parametrize null momenta

$$p_{\alpha\dot{\alpha}} = \kappa_{\alpha} \tilde{\kappa}_{\dot{\alpha}} = (1, z)^{\top} \omega (1 \ \bar{z}). \quad (1)$$

Single particle state labelled by $|\omega, z, \bar{z}, s\rangle$, diagonalize Lorentz boost by Mellin transform

$$\int_0^{\infty} \frac{d\omega}{\omega} \omega^{\Delta} |\omega, z, \bar{z}, s\rangle = |\Delta, z, \bar{z}, s\rangle. \quad (2)$$

Quick background on celestial holography

The *dynamical* statement is:

Scattering amplitudes $A(p_1(z_1, \omega_1), \dots, p_n(z_n, \omega_n))$ in 4d

\iff Correlation function $\tilde{A}(\mathcal{O}_1(z_1, \Delta_1), \dots, \mathcal{O}_n(z_n, \Delta_n))$ in 2d

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In (2,2)-signature, the OPE limit $z_i \rightarrow z_j$ in 2d implies

$p_i \cdot p_j = \omega_i \omega_j z_{ij} \bar{z}_{ij} \rightarrow 0$ in 4d.

Holomorphic collinear limit of scattering amplitudes in 4d

\iff Celestial OPE limit on 2d celestial sphere

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Does $\lim_{z_i \rightarrow z_j} \mathcal{O}_i(z_i) \mathcal{O}_j(z_j)$ have an OPE like structure?

Yes! Simply by Mellin transforming momentum space collinear splitting functions.

Example OPE

Celestial OPE between two positive helicity gluons:

$$\mathcal{O}_{\Delta_i}^a(z_i, \bar{z}_i) \mathcal{O}_{\Delta_j}^b(z_j, \bar{z}_j) \sim \frac{f^{ab}_c}{z_{ij}} B(\Delta_i - 1, \Delta_j - 1) \mathcal{O}_{\Delta_i + \Delta_j - 1}^c(z_j, \bar{z}_j) \quad (3)$$

where $B(x, y) = \frac{\Gamma(x)\Gamma(y)}{\Gamma(x+y)}$ is the Euler Beta function.

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where $B(x, y) = \frac{\Gamma(x)\Gamma(y)}{\Gamma(x+y)}$ is the Euler Beta function.

Including all the $SL(2, \mathbb{R})_R$ or \bar{z}_{ij} descendants:

$$\mathcal{O}_{\Delta_i}^a(z_i, \bar{z}_i) \mathcal{O}_{\Delta_j}^b(z_j, \bar{z}_j) \sim \frac{f^{ab}_c}{z_{ij}} \sum_{m=0}^{\infty} B(\Delta_i + m - 1, \Delta_j - 1) \frac{\bar{z}_{ij}^m}{m!} \bar{\partial}_j^m \mathcal{O}_{\Delta_i + \Delta_j - 1}^c(z_j, \bar{z}_j) \quad (4)$$

[Fan-Fotopoulos-Taylor '19] [Pate-Raclariu-Strominger-Yuan '19]

[Himwich-Pate-Singh '21] [Guevara-Himwich-Pate-Strominger '21] . . .

Question: Is there a dynamical theory generating such OPEs without referring to the symmetry or kinematical constraints in the theory?

Twistor string worldsheet theory

The Berkovits-Witten twistor string theory is a chiral worldsheet model governing degree d holomorphic maps $Z^I(\sigma)$ from a Riemann sphere Σ to supertwistor space $\mathbb{PT} \cong \mathbb{CP}^{3|4} \setminus \mathbb{CP}^1$ with

$$Z^I = (Z^A, \chi^a) = (\mu^{\dot{\alpha}}, \lambda_{\alpha}, \chi^a), \quad \alpha = 0, 1, \dot{\alpha} = \dot{0}, \dot{1}, \quad a = 1, \dots, 4.$$

Vertex operator of + helicity gluon:

$$\mathcal{O}^a = \int_{\Sigma} d\sigma j^a(\sigma) \int_{\mathbb{C}^*} \frac{ds}{s} \bar{\delta}^2(\kappa - s \lambda(\sigma)) e^{is[\mu(\sigma) \tilde{\kappa}]}. \quad (5)$$

j^a from current algebra system on the worldsheet, Kac-Moody algebra formally with level $k \rightarrow 0$. Relevant OPE:

$$j^a(\sigma_i) j^b(\sigma_j) \sim \frac{f^{ab}_c j^c(\sigma_j)}{\sigma_{ij}}. \quad (6)$$

[Witten '03] [Berkovits '04]

Celestial OPEs from the worldsheet (leading order)

Simply taking raw OPE between two + helicity gluons gives:

$$\mathcal{O}^a(z_i, \bar{z}_i) \mathcal{O}^b(z_j, \bar{z}_j) \sim \int d\sigma_i \frac{f^{ab}_c j^c(\sigma_j)}{\sigma_{ij}} \frac{ds_i}{s_i} \frac{ds_j}{s_j} \bar{\delta}^2(z_i - s_i \lambda(\sigma_i)) \bar{\delta}^2(z_j - s_j \lambda(\sigma_j)) \exp\left(i\omega_i s_i [\mu(\sigma_i) \bar{z}_i] + i\omega_j s_j [\mu(\sigma_j) \bar{z}_j]\right). \quad (7)$$

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Integrating w.r.t one of the holomorphic delta functions and integrate by parts to localize on $\frac{1}{\sigma_{ij}}$:

$$\mathcal{O}^a(z_i, \bar{z}_i) \mathcal{O}^b(z_j, \bar{z}_j) \sim f^{ab}{}_c \int d\sigma_i j^c(\sigma_j) \bar{\delta}(\sigma_{ij}) \frac{ds_j}{s_j} \frac{\langle \iota \lambda(\sigma_i) \rangle}{\langle z_i \lambda(\sigma_i) \rangle} \bar{\delta}^2(z_j - s_j \lambda(\sigma_j)) \exp\left(i\omega_i \frac{[\mu(\sigma_i) \bar{z}_i]}{\langle \iota \lambda(\sigma_i) \rangle} + i\omega_j s_j [\mu(\sigma_j) \bar{z}_j]\right). \quad (8)$$

with ι_α some reference spinor.

Celestial OPEs from the worldsheet (leading order)

Delta functions automatically localise on worldsheet OPE limit $\sigma_i \rightarrow \sigma_j$, which correspond to OPE limit on celestial sphere $z_i \rightarrow z_j$.

At leading order in z_{ij} :

Worldsheet OPE limit on twistor string worldsheet

\iff Celestial OPE limit on 2d celestial sphere

After some manipulations

\implies The celestial OPE at leading order emerges!

Provides a dynamical way to generate celestial OPEs at leading order.

All orders in z_{ij} ?

The putative CCFT is not unitary nor local

⇒ Need full OPE data

- Previous work uses Virasoro and Kac-Moody descendants to organize order by order in z_{ij} in the MHV sector
- Here instead use soft gluon descendants

[Banerjee-Ghosh-Gonzo '20] [Banerjee-Ghosh-Paul '20] [Ebert-Sharma-Wang '21]

[Banerjee-Ghosh '20] [Banerjee-Ghosh-Samal '21] [Banerjee-Ghosh-Paul '21]

Worksheet at MHV level

At MHV level, degree 1 holomorphic maps $Z^A(\sigma) = U_\alpha^A \sigma^\alpha$, fix $GL(2, \mathbb{C})$ with $\lambda_\alpha(\sigma) = \sigma_\alpha$ and remaining moduli $\mu^{\dot{\alpha}} = x^{\alpha\dot{\alpha}} \lambda_\alpha$.

Doing the integral w.r.t the delta function in positive helicity gluon operator gives operator on celestial sphere in an effective 2d CFT:

$$\mathcal{O}^a(z, \tilde{\kappa}) = \int_{\Sigma} d\lambda j^a(\lambda) \bar{\delta}(\lambda_\alpha - \kappa_\alpha) e^{i[\mu(\lambda)\tilde{\kappa}]} = j^a(z) e^{i[\mu(z)\tilde{\kappa}]}, \quad (9)$$

where we have only parametrized $\kappa_\alpha = (1 \ z)$ using 2d coordinates on the celestial sphere.

worksheet $\Sigma \xleftrightarrow{\text{holomorphically}}$ celestial sphere

Worksheet at MHV level

The effective 2d CFT on the celestial sphere:

$$S^{\text{MHV}} = \frac{1}{2\pi} \int_{S^2} \tilde{\lambda}_{\dot{\alpha}}(z) \bar{\partial} \mu^{\dot{\alpha}}(z) + S_{\text{current}} \Big|_{S^2} . \quad (10)$$

To capture all orders OPE between vertex operators, we also look at all orders OPE in the worksheet CFT:

$$j^a(z_i) j^b(z_j) = \frac{f^{ab}{}_c j^c(z_j)}{z_{ij}} + \sum_{n=0}^{\infty} \frac{z_{ij}^n}{n!} \underbrace{: \partial^n j^a j^b :}_{\frac{1}{2\pi i} \oint_{C(z_j)} dz_i \frac{\partial^n j^a(z_i) j^b(z_j)}{z_{ij}}} . \quad (11)$$

Soft gluon vertex operator

For the positive helicity gluon operator, we can expand in powers of the energy scale ω

$$\mathcal{O}^a(z, \omega(1 \bar{z})) = \sum_k \omega^k \frac{i^k}{k!} j^a(z) [\mu(z) \bar{z}]^k = \sum_k \frac{i^k}{k!} j^a(z) [\mu(z) \tilde{\kappa}]^k. \quad (12)$$

And define soft gluon vertex operators as

$$J^a[k](z, \tilde{\kappa}) := \frac{i^k}{k!} j^a(z) [\mu(z) \tilde{\kappa}]^k = \sum_{n+l=k} \frac{\tilde{\kappa}_0^n \tilde{\kappa}_1^l}{n! l!} J^a[n, l], \quad k \in \mathbb{Z}_{\geq 0}, \quad (13)$$

with $J^a[n, l] = i^{n+l} j^a(z) (\mu^0)^n(z) (\mu^1)^l(z)$.

All orders collinear expansion in momentum space

Computing the ++ gluon OPE gives

$$\begin{aligned} \mathcal{O}_i^a(z_i, \tilde{\kappa}_i) \mathcal{O}_j^b(z_j, \tilde{\kappa}_j) &= j^a(z_i) j^b(z_j) e^{i[\mu(z_i) \tilde{\kappa}_i]} e^{i[\mu(z_j) \tilde{\kappa}_j]} \\ &= \left(\sum_{m=0}^{\infty} z_{ij}^{m-1} j_{-m}^a j^b \right) \left(\sum_{n=0}^{\infty} \frac{z_{ij}^n}{n!} \partial_{z_j}^n \left(e^{i[\mu \tilde{\kappa}_i]} \right) e^{i[\mu \tilde{\kappa}_j]} \right), \end{aligned} \quad (14)$$

normal ordered and evaluated at z_j where

$$j_{-m}^a j^b := \begin{cases} f^{ab}{}_c j^c(z_j) & \text{when } m = 0; \\ \frac{\partial_{z_j}^{m-1} j^a j^b(z_j)}{(m-1)!} & \text{when } m \geq 1. \end{cases} \quad (15)$$

All orders collinear expansion in momentum space

We choose to organize the OPEs using soft gluon operators, their action on $\mathcal{O}^a(z, \tilde{\kappa})$:

$$\begin{aligned} & \mathcal{J}^a[k](z_i, \tilde{\kappa}_i) \mathcal{O}_j^b(z_j, \tilde{\kappa}_j) \\ &= \frac{[i \partial_j]^k f^{ab}_c \mathcal{O}_j^c(z_j, \tilde{\kappa}_j)}{k! z_{ij}} + \sum_{p=1}^{\infty} z_{ij}^{p-1} \underbrace{\frac{i^k}{k!} \sum_{n=0}^p \frac{1}{n!} j_{n-p}^a j^b \partial_{z_j}^n \left([\mu \tilde{\kappa}_i]^k \right) e^{i[\mu \tilde{\kappa}_j]}}_{\mathcal{J}_{-p}^a[k](\tilde{\kappa}_i) \mathcal{O}_j^b(z_j, \tilde{\kappa}_j)} . \end{aligned}$$

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$$\begin{aligned} & J^a[k](z_i, \tilde{\kappa}_i) \mathcal{O}_j^b(z_j, \tilde{\kappa}_j) \\ &= \frac{[i \partial_j]^k}{k!} \frac{f^{ab}_c \mathcal{O}_j^c(z_j, \tilde{\kappa}_j)}{z_{ij}} + \underbrace{\sum_{p=1}^{\infty} z_{ij}^{p-1} \frac{i^k}{k!} \sum_{n=0}^p \frac{1}{n!} J_{n-p}^a j^b \partial_{z_j}^n \left([\mu \tilde{\kappa}_i]^k \right)}_{J_{-p}^a[k](\tilde{\kappa}_i) \mathcal{O}_j^b(z_j, \tilde{\kappa}_j)} e^{i[\mu \tilde{\kappa}_j]} . \end{aligned}$$

Reorganizing the OPE (14) using the soft descendants $J_{-p}^a[k]$:

$$\mathcal{O}_i^a \mathcal{O}_j^b = \sum_{p=0}^{\infty} z_{ij}^{p-1} \sum_{k=0}^p \sum_{\ell=0}^{p-k} \frac{(-[\tilde{\kappa}_i \partial_{\tilde{\kappa}_j}])^\ell}{\ell!} J_{-p}^a[k](\tilde{\kappa}_i) \mathcal{O}^b(z_j, \tilde{\kappa}_i + \tilde{\kappa}_j) . \quad (16)$$

This is the all orders collinear expansion in the MHV sector.

All orders $SL(2, \mathbb{R})_L$ descendants in conformal primary basis

We can also parametrize $\tilde{\kappa}_{\dot{\alpha}} = \omega\left(\frac{1}{\bar{z}}\right)$ and Mellin transform in ω to obtain the all-order celestial OPE in conformal primary basis:

$$\begin{aligned} \mathcal{O}_{+,\Delta_i}^a(z_i, \bar{z}_i) \mathcal{O}_{+,\Delta_j}^b(z_j, \bar{z}_j) &= \sum_{p=0}^{\infty} z_{ij}^{p-1} \sum_{k=0}^p \sum_{\ell=0}^{p-k} \sum_{m=0}^{\infty} \frac{\bar{z}_{ij}^m}{m!} \frac{\bar{D}_k^\ell}{\ell!} J_{-p}^a[k](\bar{z}_i) \\ &\times B(\Delta_i + k + \ell + m - 1, \Delta_j - 1) \bar{\partial}_j^m \mathcal{O}_{+,\Delta_i+\Delta_j+k-1}^b(z_j, \bar{z}_j). \quad (17) \end{aligned}$$

with $\bar{D}_j = -\bar{z}_{ij} \bar{\partial}_j + \Delta_i + \Delta_j + k - 3$.

Discussion

- This gives a genuinely dynamical CFT on 2d celestial sphere to generate observables to match those from gauge theory in 4d, which provides a good starting point for a bottom up construction of celestial holography
- The momentum space formulae gives all-orders collinear expansion and the one in conformal primary basis gives all-orders celestial OPE
- Can obtain all-order expressions for other helicity and orientation configurations
- Matching with known results in low orders in z_{ij} using null state
- Relation to infinite dimensional chiral algebra living on the celestial sphere
- Beyond MHV sector?