

Non-perturbative Double Copy in Flatland

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with

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QCD meets Gravity

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Motivation

Double-copy (DC) – one of our best tools.

Want to do it off-shell.

Selling points

1. Easy: Would make DC trivial
2. Explain: Would help explain why DC works
3. New: Could DC new, off-shell quantities

Theories with some amount of manifest color-kinematics duality:

- Self-dual Yang-Mills (SDYM) [Monteiro & O'Connell '11] – tree level, vanishes on-shell
- 3D Chern-Simons theory [Ben-Shahar & Johansson '21] – loop level, topological
- Non-abelian fluid [Cheung & JM '20] – tree level
- Non-linear sigma model (NLSM) [Cheung & Shen '16; Cheung & JM '21] – tree level

Still more to learn

This talk

Pion-like theory in 2D that is

- interacting
- manifest color-kinematics duality
- all loop orders/off-shell/Lagrangian-level

NLSM review

Chiral current j_μ^a

$$(1) F_{\mu\nu}^c(j) = \partial_{[\mu} j_{\nu]}^c + f_{ab}{}^c j_\mu^a j_\nu^b = 0$$

Pure gauge, implicitly defines pion, field basis
indep $j_\mu = ig^{-1} \partial_\mu g$

$$(2) \partial^\mu j_\mu^a = 0$$

Chiral current is conserved, Lorenz gauge, pion
EOM $\partial(g^{-1} \partial g) = 0$

Cubic EOM with *manifest* BCJ

Combine (1) $F = 0$ and (2) $\partial j = 0$

$$\square j_{\mu}^c + f_{ab}{}^c j^{a\nu} \partial_{\nu} j_{\mu}^b = 0$$

Manifestly satisfies kinematic Jacobi *off-shell* at tree level. Related construction for YM. Yields formula for numerators for NLSM & YM
[Cheung & JM '21].

Two dimensions

Easy to satisfy (2) $\partial_\mu j^\mu = 0$ in 2D by dualizing chiral current $j^\mu = \epsilon^{\mu\nu} \partial_\nu \phi = \tilde{\partial}^\mu \phi$ where $\tilde{\partial}^\mu = \epsilon^{\mu\nu} \partial_\nu$.

$$\square \phi^c - \frac{1}{2} f_{ab}{}^c \partial_\mu \phi^a \tilde{\partial}^\mu \phi^b = 0$$

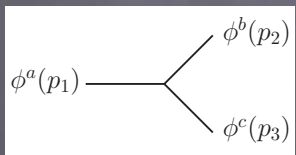
Unlike NLSM, this EOM can be *integrated* to get the Lagrangian for Zakharov-Mikhailov (ZM) theory

$$\mathcal{L}_{\text{ZM}} = \frac{1}{2} \partial_\mu \phi_a \partial^\mu \phi^a + \frac{1}{3!} f_{abc} \phi^a \partial_\mu \phi^b \tilde{\partial}^\mu \phi^c$$

Feynman rule

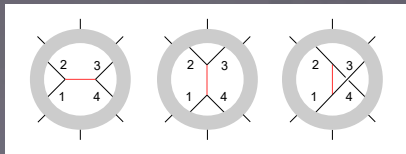
$$\mathcal{L}_{\text{ZM}} = \frac{1}{2} \partial_\mu \phi_a \partial^\mu \phi^a + \frac{1}{3!} f_{abc} \phi^a \partial_\mu \phi^b \tilde{\delta}^\mu \phi^c$$

Define $\langle ij \rangle = \epsilon_{\mu\nu} p_i^\mu p_j^\nu$


$$= -if_{abc} \langle 12 \rangle$$

Secretly Bose sym off-shell $\langle 12 \rangle = \langle 23 \rangle = \langle 31 \rangle$

Key step: check off-shell Jacobi



$$\frac{c_s n_s}{s} + \frac{c_t n_t}{t} + \frac{c_u n_u}{u}$$

Color-stripped vertex $\sim \langle ij \rangle = \epsilon_{\mu\nu} p_i^\mu p_j^\nu$

Off-shell Jacobi

$$n_s + n_t + n_u = \langle 12 \rangle \langle 34 \rangle + \text{cyclic}(1, 2, 3) = 0$$

Just Schouten identity.

ZM manifest BCJ to all loop orders.

Selling point (1): Easy

Read off replacement rules from \mathcal{L}_{ZM}

$$\begin{aligned} V^a &\rightarrow V \\ f_{ab}{}^c V^a W^b &\rightarrow \partial_\mu V \tilde{\partial}^\mu W \\ g_{ab} V^a W^b &\rightarrow \int VW. \end{aligned}$$

Double copy literally reduced to a replacement.

$$\mathcal{L}_{\text{BAS}} \rightarrow \mathcal{L}_{\text{ZM}} \rightarrow \mathcal{L}_{\text{SG}}$$

Double-copy replacement

$$f_{ab}{}^c V^a W^b \rightarrow \partial_\mu V \tilde{\partial}^\mu W$$

$$\mathcal{L}_{\text{BAS}} = \frac{1}{2} \partial_\mu \phi_{a\bar{a}} \partial^\mu \phi^{a\bar{a}} + \frac{1}{3!} f_{abc} f_{\bar{a}\bar{b}\bar{c}} \phi^{a\bar{a}} \phi^{b\bar{b}} \phi^{c\bar{c}}$$

↓

$$\mathcal{L}_{\text{ZM}} = \frac{1}{2} \partial_\mu \phi_a \partial^\mu \phi^a + \frac{1}{3!} f_{abc} \phi^a \partial_\mu \phi^b \tilde{\partial}^\mu \phi^c$$

↓

$$\mathcal{L}_{\text{SG}} = \frac{1}{2} \partial_\mu \phi \partial^\mu \phi + \frac{1}{3!} \phi \partial_\mu \partial_\nu \phi \tilde{\partial}^\mu \tilde{\partial}^\nu \phi$$

Selling point (2): Explains

Kinematica algebra is area-preserving diffeomorphisms on torus (Poisson). Isomorphic to large N limit of color algebra $U(N)$ [Hoppe '88].

$$\lim_{N \rightarrow \infty} U(N) \sim \text{Diff}_{S^1 \times S^1}$$

Selling point (3): new
interesting things

Masses & Higher Dim Ops

Masses

Masses are fine, only affect kinetic term

Higher dim ops

ZM^2 = special Galileons are secretly free in 2D
[de Rham, Fasiello & Tolley '14].

No worries, infinitely many higher dim ops
double-copy

Higher Dim Ops Continued

Can double-copy any operator except

1. No closed color loops
2. No multiple color traces

because $g_{ab} V^a W^b \rightarrow \int VW$.

Example ϕ^4 :

$$\mathcal{O}_{\text{BAS}} = f_{abe} f_{cde} f_{\bar{a}\bar{b}\bar{e}} f_{\bar{c}\bar{d}\bar{e}} \phi^{a\bar{a}} \phi^{b\bar{b}} \phi^{c\bar{c}} \phi^{d\bar{d}}$$

$$\mathcal{O}_{\text{ZM}} = f_{abe} f_{cde} \partial_\mu \phi^a \tilde{\partial}^\mu \phi^b \partial_\nu \phi^c \tilde{\partial}^\nu \phi^d$$

$$\mathcal{O}_{\text{SG}} = \partial_\mu \partial_\nu \phi \tilde{\partial}^\mu \tilde{\partial}^\nu \phi \partial_\rho \partial_\sigma \phi \tilde{\partial}^\rho \tilde{\partial}^\sigma \phi.$$

Moyal deformation

Moyal deformation, also in SDYM [Chacón, García-Compeán, Luna, Monteiro & White '14]

$$f_{ab}{}^c V^a W^b \rightarrow \partial_\mu V \tilde{\partial}^\mu W \quad (\text{above})$$

$$f_{ab}{}^c V^a W^b \rightarrow \frac{N}{2\pi} \sin\left(\frac{2\pi}{N} \partial_V \tilde{\partial}_W\right) VW \quad (\text{Moyal})$$
$$\sim \sin(\langle 12 \rangle)$$

Still satisfies Jacobi at full off-shell Lagrangian level

Integrability & Wilson line

- ZM known to be integrable
- Lax pair formed from Wilson line
- Infinite tower of off-shell currents
- Everything double-copies to SG via replacement rules

Similar story for SDYM [Chacón,
García-Compeán, Luna, Monteiro & White '14]

Details of currents 1/3

Color \rightarrow diff

$$\underline{V} = V^a \underline{T}_a \quad \rightarrow \quad \underline{V} = \partial_\mu V \tilde{\partial}^\mu$$

$$[V^a \underline{T}_a, W^b \underline{T}_b] = f_{ab}{}^c V^a W^b \underline{T}_c$$

$$\rightarrow [\partial_\mu V \tilde{\partial}^\mu, \partial_\nu W \tilde{\partial}^\nu] = \partial_\mu (\partial_\nu V \tilde{\partial}^\nu W) \tilde{\partial}^\mu$$

$$c.f. \quad f_{ab}{}^c V^a W^b \quad \rightarrow \quad \partial_\mu V \tilde{\partial}^\mu W$$

Details of currents 2/3

$$\underline{A}_\mu = \frac{1}{1 - \lambda^2} (\tilde{\partial}_\mu \underline{\Phi} + \lambda \partial_\mu \underline{\Phi})$$

$$\underline{F}_{\mu\nu} = 0$$

$$\underline{W}(x) = \text{P exp} \left[- \int^x dx'^\mu \underline{A}_\mu(x') \right]$$

$$\underline{J}_\mu = \tilde{\partial}_\mu \underline{W} = \sum_{k=0}^{\infty} \lambda^{-k} \underline{J}_\mu^{(k)}$$

Details of currents 3/3

Example ZM currents

$$\begin{aligned}\underline{J}_{\underline{\mu}}^{(1)} &= \tilde{\partial}_{\underline{\mu}}\underline{\phi}, & \underline{J}_{\underline{\mu}}^{(2)} &= \partial_{\underline{\mu}}\underline{\phi} + \tilde{\partial}_{\underline{\mu}}\underline{\phi}\underline{\phi} \\ \underline{J}_{\underline{\mu}}^{(3)} &= \tilde{\partial}_{\underline{\mu}}\underline{\phi} + \partial_{\underline{\mu}}\underline{\phi}\underline{\phi} + \tilde{\partial}_{\underline{\mu}}\underline{\phi} \int^x dx' \tilde{\partial}\underline{\phi} \\ &+ \tilde{\partial}_{\underline{\mu}}\underline{\phi} \int^x dx' \partial\underline{\phi}\underline{\phi}\end{aligned}$$

Double copy to SG $\underline{V} = V^a \underline{T}_a \rightarrow \underline{V} = \partial_{\underline{\mu}} V \tilde{\partial}^{\underline{\mu}}$

$$\underline{J}_{\underline{\mu}}^{(1)} = \tilde{\partial}_{\underline{\mu}} \partial_{\underline{\nu}} \phi \tilde{\partial}^{\underline{\nu}}, \quad \underline{J}_{\underline{\mu}}^{(2)} = \partial_{\underline{\mu}} \partial_{\underline{\nu}} \phi \tilde{\partial}^{\underline{\nu}} + \tilde{\partial}_{\underline{\mu}} \partial_{\underline{\nu}} \phi \tilde{\partial}^{\underline{\nu}} \partial_{\underline{\rho}} \phi \tilde{\partial}^{\underline{\rho}}$$

Can do the whole tower (closed form)

Numerical double-copy!

$$\phi_{\text{ZM}}^a(x) T^a = \int d^2p e^{ipx} T_p \phi_{\text{SG}}(p)$$

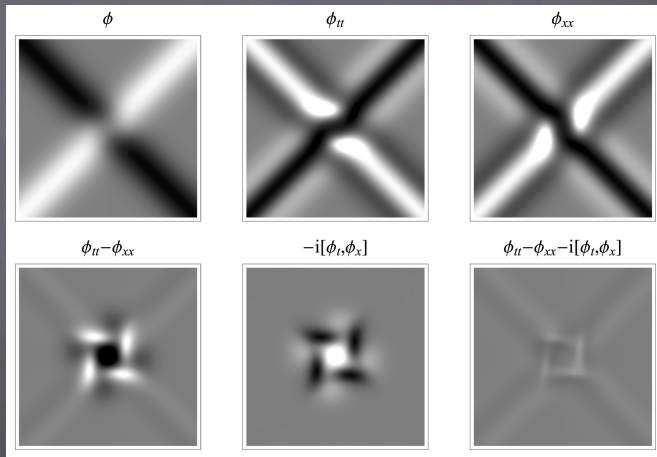
Discretize, take $N = 499$ large

Double-copy from SG to ZM (backwards)

KLT: $A_{\text{GR}} \sim s^{n-3} A_{\text{YM}} A_{\text{YM}} + \dots$

Collide two SG Gaussian wave packets.
Plot ZM sol projected onto $\sum_p T_p$.

Collide two SG Gaussian wave packets.
Plot ZM sol projected onto $\sum_p T_p$.



Overview

- Manifest color-kinematics duality for 2D “pion” at Lagrangian level
- Double copy is just replacement
- Kinematic algebra is diff algebra
- Color-kinematics duality – Hoppe
- Double copy masses, higher dim ops, Wilson lines, numerical solutions...

Thank you!