

Kinematic Hopf algebra(KiHA) in amplitude and form factor



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Brandhuber, GC, Travaglini, Wen(2104.11206, 2108.04216)

Brandhuber, GC, Johansson, Travaglini, Wen(2111.15649)

Brandhuber, Brown,GC,Gowdy, Travaglini, Wen(2208.05886)

GC, Lin, Wen(2208.05519)

Bjerrum-Bohr, GC, Skowronek (2301.?????)

QCD Meet Gravity @Zürich

Contents

- Background
- Framework: Kinematic Hopf algebra(KiHA)
- Double and applications to (spinning) heavy-mass effective theory

Motivation

- Develop a systematic framework to construct the amplitude and form factor
- Exploring the underlying algebra structures in colour-kinematic and double copy duality
- Constructing new theory(effective theory & classical/high spin theory)

Gauged Bi-adjoint scalar theory

Heavy matter coupling to gluon (Yang-Mills)

$$\mathcal{L} = \mathcal{L}^{\text{YMS}} + m^2 \phi^{I,a} \phi^{I,a} + \frac{g'}{3!} f^{IJK} f^{abc} \phi^{I,a} \phi^{J,b} \phi^{K,c}$$

Motivation

- Have c-k duality
- Double copy to graviton-scalar, EYM, EM amplitude
- Closely related to gravitational scattering (including, binary black hole, Compton scattering, light bending et al)
- It's fun (Novel relations among the BCJ numerator)

Chiodaroli, Jin, Radu 2013,
Chiodaroli, Gunaydin, Hohansson, Roiban 2014,
Cachazo, He and Yuan 2014

Background on color-kinematic(CK) and double copy (DC)

Bern, Carrasco, Johansson 2008,2010

CK duality

Sum over all the cubic graphs

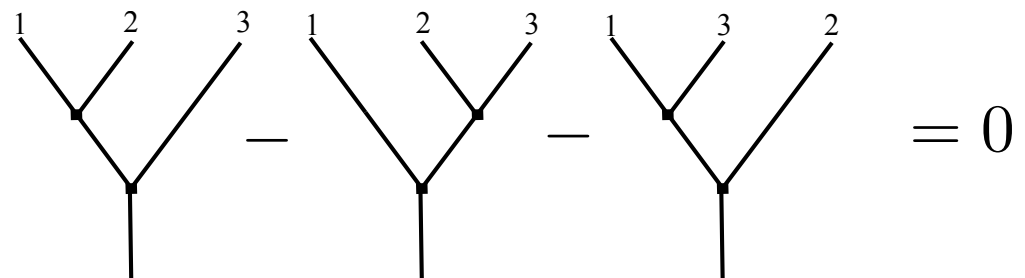
$$\mathcal{A}_n^{\text{tree}} = g^{n-2} \sum_{I=1}^{(2n-5)!!} \frac{c_I N_I}{D_I}$$

Scalar propagator

c_I Color factor

N_I kinematic numerator

$$f^{abe} f^{ecd} + f^{ace} f^{edb} + f^{ade} f^{ebc} = 0$$



DC duality

$$\mathcal{A}_n^{\text{tree}} = g^{n-2} \sum_{I=1}^{(2n-5)!!} \frac{\bar{N}_I N_I}{D_I}$$

Factor out the propagators

Kinematic Hopf algebra(KiHA) in Bi-adjoint scalar

Wen 2022, GC, Lin, Wen 2022

Building Blocks in KiHA

Brandhuber, GC, Johansson, Travaglini, Wen 2021-11

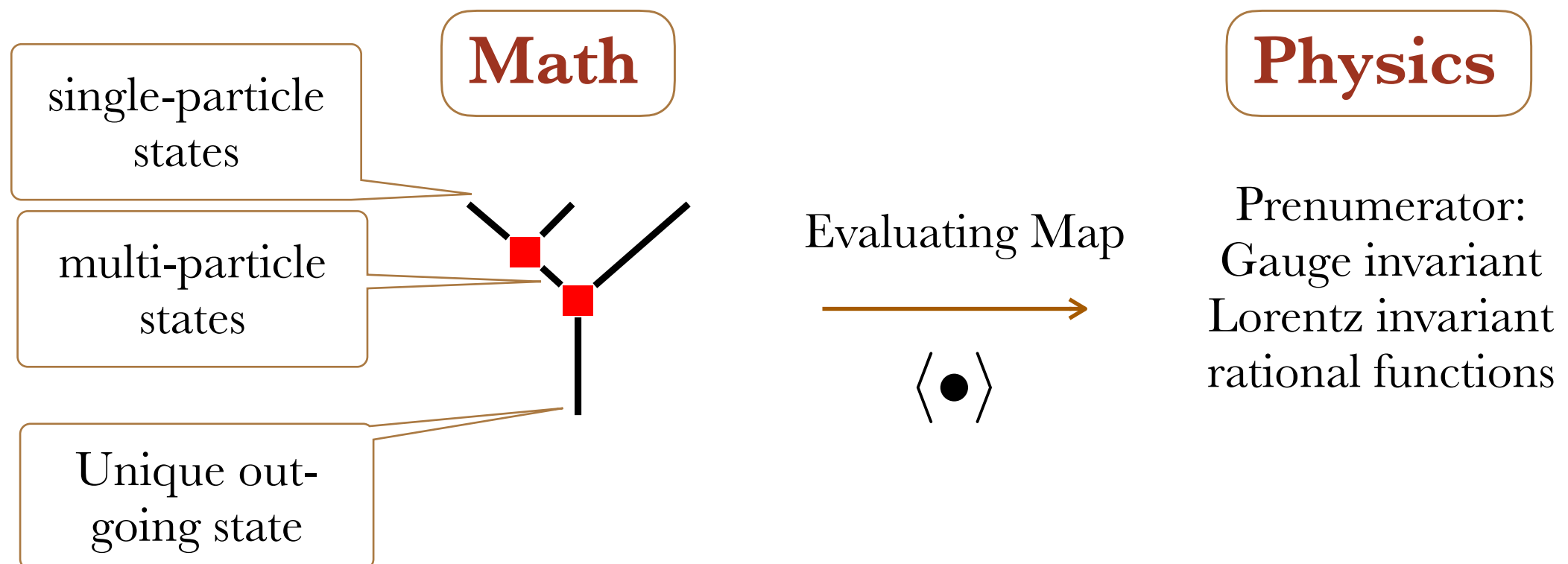
Brandhuber, Brown, GC, Gowdy, Travaglini2022;

GC,Johansson, Teng, Wang 2019; GC,Johansson,Teng,Wang 2021

1. Generators T
2. Fusion product \star
3. Evaluating Map $\langle \bullet \rangle$

Feynman Rule

- Quantum Field
- Vertex: Interaction
- Contract with external states



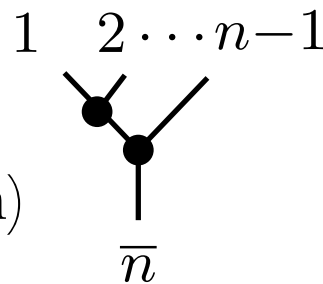
Construct amplitude

KiHA

- Defined the commutator

$$[X, Y] \equiv X \star Y - Y \star X$$

- Perform commutators in a given order (a cubic graph)

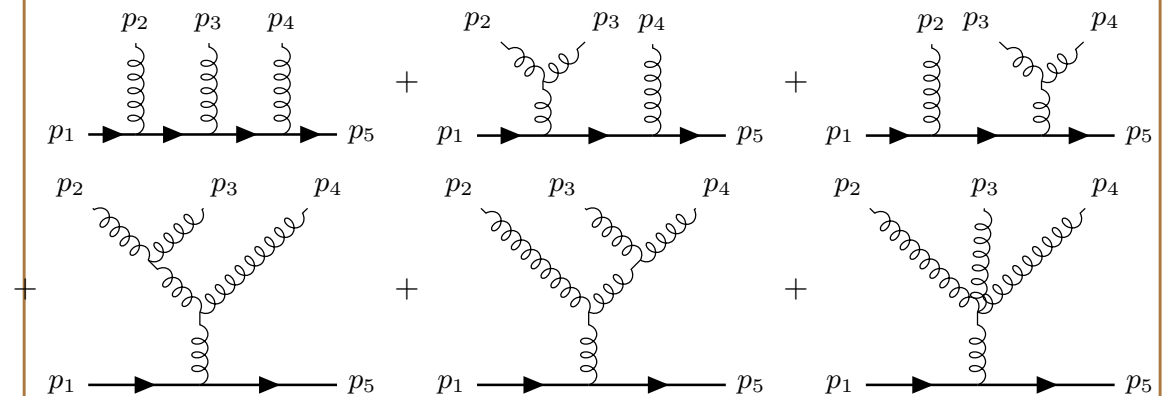


- Sum over all the cubic graph, e.g. colour ordered amplitude

$$A(\sigma, \bar{n}) = \sum_{\Gamma \in R_\sigma} \frac{\langle \hat{\mathcal{N}}(\Gamma) \rangle}{d_\Gamma} \sigma(1) \cdots \sigma(n-1)$$

Feynman Rule

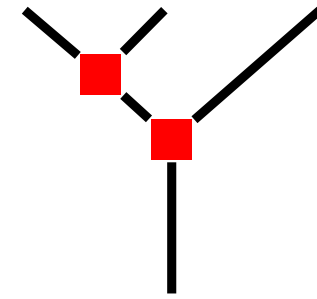
Sum over all the Feynman diagrams



**Gauge invariant
function
from each cubic graph**

Algebraic generators

$$K_i = \begin{cases} T_{(i)}^{(i)} & \text{for gluons} \\ T_{(i)}^{(i)} t^{a_i} & \text{for scalars} \\ T_{(\tau_1), \dots, (\tau_r)}^{(\alpha)} t^{a_i} \dots t^{a_j} & \text{intermediate states} \end{cases}$$



α

fusion product order

t^{a_i}

flavour group generator

$(\tau_1), \dots, (\tau_r)$

partition of gluon labels

$$T_{(12), (43), (5)}^{(6124357)} t^{a_6} t^{a_7}$$

Fusion product

Brandhuber, GC, Johansson, Travaglini, Wen 2021-11

Non-abelian quasi-shuffle Hopf algebra

$$\begin{aligned}
 & T_{(\tau_1), \dots, (\tau_r)}^{(\alpha)} \star T_{(\omega_1), \dots, (\omega_s)}^{(\beta)} \\
 &= \sum_{\substack{\pi|_{\tau} = \{(\tau_1), \dots, (\tau_r)\} \\ \pi|_{\omega} = \{(\omega_1), \dots, (\omega_s)\}}} (-1)^{t-r-s} T_{(\pi_1), \dots, (\pi_t)}^{(\alpha\beta)},
 \end{aligned}$$

where $\pi|_{\tau}$ (or $\pi|_{\omega}$) means a restriction to the elements of π in τ (or ω), e.g. $\{(235), (4), (678)\}|_{\{2,3,4,8\}} = \{(23), (4), (8)\}$.

Example

$$\begin{aligned}
 T_{(1),(2)}^{(12)} \star T_{(34)}^{(345)} &= T_{(1),(2),(34)}^{(12345)} + T_{(1),(34),(2)}^{(12345)} + T_{(34),(1),(2)}^{(12345)} \\
 &\quad - T_{(1),(234)}^{(12345)} - T_{(134),(2)}^{(12345)}
 \end{aligned}$$

Shuffle

Stuffing

Super index is critical for non-trivial commutator

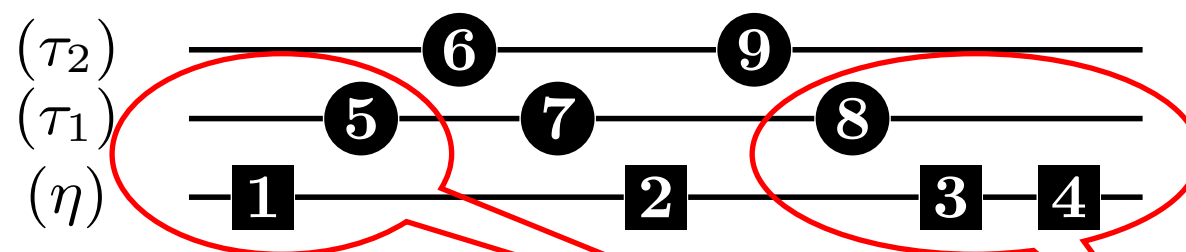
Evaluation map

From algebraic generators to gauge invariants

$$t^\eta T_{(\tau_1), \dots, (\tau_r)}^{(\alpha)} \xrightarrow[\text{map}]{\langle \bullet \rangle} \text{tr}(t^\eta t^{a_n}) \langle T_{(\tau_1), \dots, (\tau_r)}^{(\alpha)} \rangle_m$$

$$\langle T_{(\tau_1), \dots, (\tau_r)}^{(\alpha)} \rangle_m = \left(\begin{array}{c} \eta \\ \phi_1 \\ \vdots \\ \phi_{k-1} \end{array} \begin{array}{c} \tau_1 \quad \dots \quad \tau_r \\ \diagdown \quad \diagup \quad \dots \quad \diagdown \quad \diagup \\ \bullet \text{---} \bullet \text{---} \bullet \text{---} \bullet \\ \diagup \quad \diagdown \quad \dots \quad \diagup \quad \diagdown \end{array} \right)^{(\alpha)}$$

$$= \frac{2^r \prod_{i=1}^r \left(p_{\Theta_L^{(\alpha)}(\tau_i)} \cdot F_{\tau_i} \cdot p_{\Theta_R^{(\alpha)}(\tau_i)} \right)}{(p_\eta^2 - m^2)(p_{\eta\tau_1}^2 - m^2) \cdots (p_{\eta\tau_1 \cdots \tau_{r-1}}^2 - m^2)}$$



$$\langle T_{(578), (69)}^{(156729834)} \rangle_m = \frac{4p_1 \cdot F_{578} \cdot p_{34} \quad p_{15} \cdot F_{69} \cdot p_{348}}{(p_{1234}^2 - m^2)(p_{1234578}^2 - m^2)} \cdot \quad \underline{F_i^{\mu\nu} = p_i^\mu \varepsilon_i^\nu - \varepsilon_i^\mu p_i^\nu}$$

Warm up example

$$\mathcal{A}(\bar{1}, \bar{2}, 3, \bar{4}) = \begin{array}{c} \bar{1} \quad \bar{2} \quad 3 \\ \diagdown \quad \diagup \\ \bullet \\ \diagup \quad \diagdown \\ \bullet \\ | \\ \bar{4} \end{array} + \boxed{\begin{array}{c} \bar{1} \quad \bar{2} \quad 3 \\ \diagdown \quad \diagup \\ \bullet \\ \diagup \quad \diagdown \\ \bullet \\ | \\ \bar{4} \end{array}}$$

1,2,4 is scalar
3 is gluon

Example KiHA

1. External state:

$$K_i = \begin{cases} T^{(1)} t^{a_1}, & i = 1 \text{ scalar} \\ T^{(2)} t^{a_2}, & i = 2 \text{ scalar} \\ T_{(3)}^{(3)}, & i = 3 \text{ gluon} \end{cases}$$

2. Perform the commutators

$$\begin{aligned} \langle [[K_1, [K_2, K_3]]] \rangle &= \langle [T^{(1)} t^{a_1}, (T_{(3)}^{(23)} - T_{(3)}^{(32)}) t^{a_2}] \rangle \\ &= \langle (T_{(3)}^{(123)} - T_{(3)}^{(132)}) t^{a_1} t^{a_2} - (T_{(3)}^{(231)} - T_{(3)}^{(321)}) t^{a_2} t^{a_1} \rangle \end{aligned}$$

3. Sum over the cubic graphs and mapping to kinematics

$$\mathcal{A}(\bar{1}, \bar{2}, 3, \bar{4}) = \frac{2p_2 \cdot F_3 \cdot p_1}{(p_{12}^2 - m^2)(p_{23}^2 - m^2)} f^{a_1 a_2 a_4}$$

Form factor

Motivations

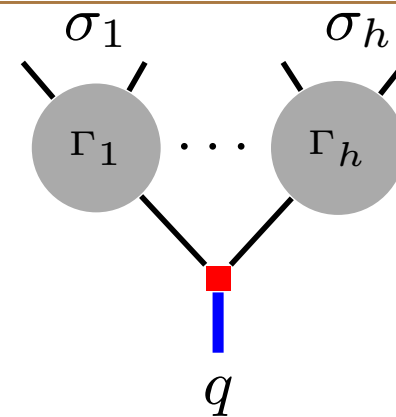
1. Does such a fusion rule have the universal behaviour as local interaction?

The operator

$$\text{Tr}(\phi^h) = \text{tr}(t_c^{A_1} \cdots t_c^{A_h}) \text{tr}(t^{a_1} \cdots t^{a_h}) \phi^{a_1, A_1} \cdots \phi^{a_h, A_h}$$

Form factor

$$\mathcal{F}_{\text{Tr}(\phi^h)}(\sigma) = \sum_{\Gamma \in \mathcal{R}_\sigma^{(h)}} \frac{\langle \hat{\mathcal{N}}(\Gamma_1) \star \cdots \star \hat{\mathcal{N}}(\Gamma_h) \rangle}{d_{\Gamma_1} \cdots d_{\Gamma_h}}$$



Fusion product (Local)

Same as amplitude

Evaluation map(global)

$$t^\eta T_{(\tau_1), \dots, (\tau_r)}^{(\alpha)} \xrightarrow[\text{map}]{\langle \bullet \rangle} \text{tr}(t^\eta) \langle T_{(\tau_1), \dots, (\tau_r)}^{(\alpha)} \rangle_q$$

$$\langle T_{(\tau_1), \dots, (\tau_r)}^{(\alpha)} \rangle_q = \langle T_{(\tau_1), \dots, (\tau_r)}^{(\alpha)} \rangle_m \Big|_{m^2 \rightarrow q^2}$$

New relation(Beyond Jacobi)

Beyond Jacobi

$$\sum_{i=1}^h \langle \widehat{\mathcal{N}}(\Gamma_1) \star \dots \star \mathcal{N}([\Gamma_i, j]) \star \dots \star \widehat{\mathcal{N}}(\Gamma_h) \rangle$$

$$= \sum_{i=1}^h \langle \text{Diagram} \rangle = 0.$$

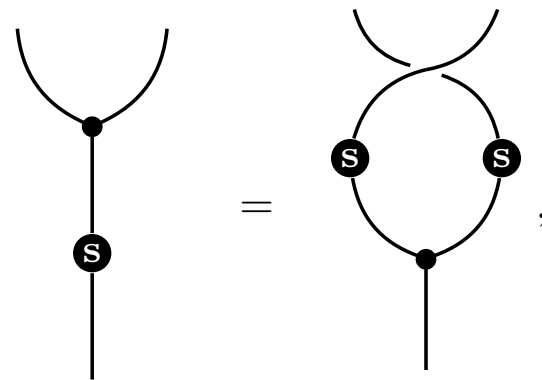
Antipode relation

$$\langle \widehat{\mathcal{N}}(12 \dots n-1) \rangle \Big|_{t^a \rightarrow \mathbb{I}} = \langle S(\widehat{\mathcal{N}}(12 \dots n-1)) \rangle \Big|_{t^a \rightarrow \mathbb{I}},$$

S is anti-morphism

$$S(X \star Y) = S(Y) \star S(X)$$

$$S(T^{(i)}) = T^{(i)}, S(T_{(j)}^{(j)}) = -T_{(j)}^{(j)}$$



More Hopf element

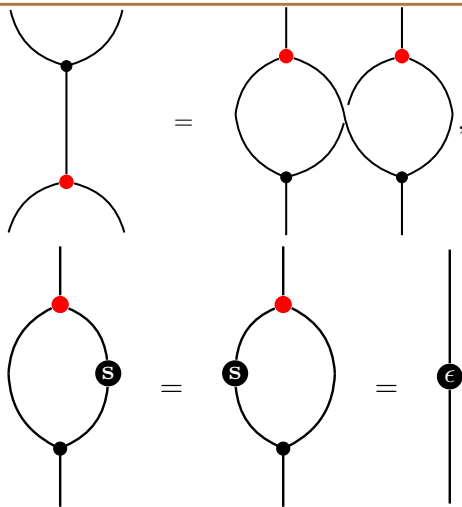
Brandhuber, Brown, GC, Gowdy, Travaglini2022;

Coproduct and counit

$$\Delta(T^{(\alpha)}) = \left(T^{(\alpha)} \otimes T^{(\alpha)} \right) ,$$

$$\Delta(T_{(\tau_1), \dots, (\tau_r)}^{(\alpha)}) = \sum_{i=0}^r \left(T_{(\tau_1), \dots, (\tau_i)}^{(\alpha)} \otimes T_{(\tau_{i+1}), \dots, (\tau_r)}^{(\alpha)} \right) ,$$

$$\epsilon(T_{\tau}^{(\alpha)}) = 0 , \epsilon(T^{(\alpha)}) = \mathbb{I}_{\text{qs}}$$



$$\Delta(T_{\tau}^{(\alpha)}) \star \Delta(T_{\omega}^{(\beta)}) = \Delta(T_{\tau}^{(\alpha)} \star T_{\omega}^{(\beta)}) ,$$

$$\epsilon(T_{\tau}^{(\alpha)}) \star \epsilon(T_{\omega}^{(\beta)}) = \epsilon(T_{\tau}^{(\alpha)} \star T_{\omega}^{(\beta)}) ,$$

$$\star(\mathbb{I}_{\text{qs}} \otimes S)\Delta(T_{\tau}^{(\alpha)}) = \star(S \otimes \mathbb{I}_{\text{qs}})\Delta(T_{\tau}^{(\alpha)}) = \epsilon(T_{\tau}^{(\alpha)})\mathbb{I}_{\text{qs}}$$

Kinematic algebra is a non-abelian extended quasi-shuffle Hopf algebra

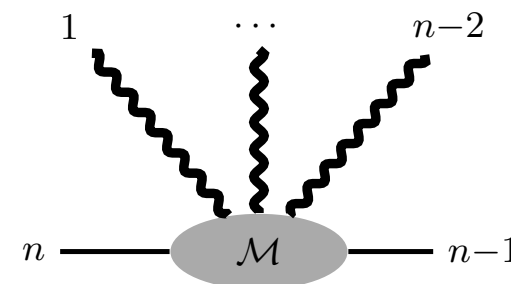
Applications: (spinning) heavy-mass effective
field theory

Double copy

Brandhuber, GC, Travaglini, Wen 2021-04;
 Brandhuber, GC, Johansson, Travaglini, Wen 2021-11
 Brandhuber, Brown, GC, Gowdy, Travaglini 2022;

Full scalar

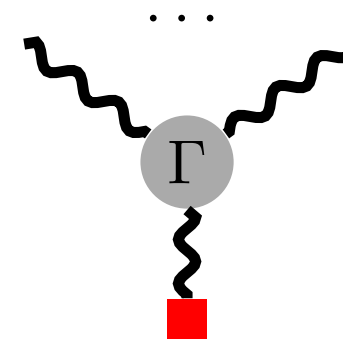
$$M(12 \dots n-2, \overline{n-1}, \overline{n}) = - \sum_{\Gamma \in \tilde{\rho}} \frac{\mathcal{N}(\Gamma, \overline{n-1}, \overline{n}) \mathcal{N}(\Gamma, \overline{n}, \overline{n-1})}{d_{\Gamma}},$$



Heavy mass limit

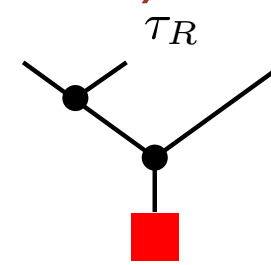
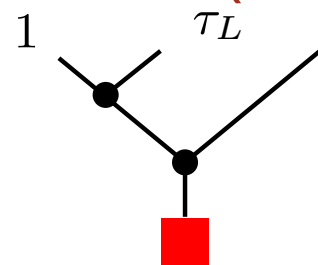
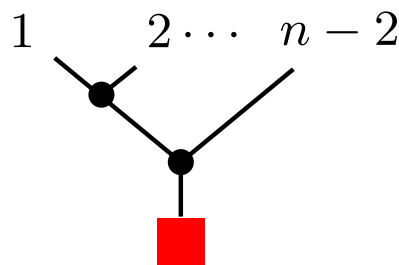
HEFT-GR

$$M(12 \dots n-2, v) = \sum_{\Gamma \in \tilde{\rho}} \frac{[\mathcal{N}(\Gamma, v)]^2}{d_{\Gamma}}$$



Novel Factorization behaviour (coproduct)

$$\mathcal{N}([12 \dots n-2], v) \xrightarrow{v \cdot p_{1\tau_L}} p_{\Theta(\tau_R)} \cdot p_{\tau_R[1]} \mathcal{N}([1\tau_L], v) \mathcal{N}([\tau_R], v),$$

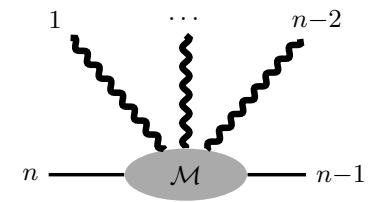


Factorisation behaviour also holds for the Compton scattering at spin infinity order

Properties of BCJ numerator

Bjerrum-Bohr, GC, Skowronek 2023

- Manifestly Gauge invariant
- Local amplitude
- Covariant(fit the PM expansion)
- Closed form for arbitrary spin order
- Unique at four-point and finite controlled freedom at a higher point.



Four point example

$$\mathcal{N}_s(12, v) = -\frac{w_1 \cdot F_1 \cdot F_2 \cdot w_2}{v \cdot p_1} + \mathcal{N}'_s(12, v).$$

$$\begin{aligned} \mathcal{N}'_s(12, v) = & (a \cdot F_2 \cdot F_1 \cdot v a \cdot p_1 - a \cdot F_2 \cdot v a \cdot F_1 \cdot p_2) G_1(x_1) G_1(x_2) \\ & + i \left((a \cdot S \cdot F_1 \cdot F_2 \cdot p_1 + a \cdot S \cdot F_2 \cdot F_1 \cdot p_2) G_2(x_1, x_2) \right. \\ & \left. + \text{tr}(S F_2 F_1) G_1(x_{12}) \right). \end{aligned}$$

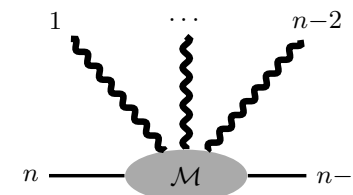
Outlook

- Geometry representation of KiHA:(ambitwistor)-string world sheet, Self-dual Yang-Mills, Chern-Simons, permutohedra

e.g. Monteiro2022; Ben-shahar, Johansson 2021;
Mafra, Schlotterer 2022; Cao, Dong, He, Zhang, Zeng 2022

- Classical/high spin kinematic algebra

e.g. Chen, Chung, Huang, Kim 2021; Bern, Luna, Roiban, Shen, Zeng 2020;
Guevara, Ochirov, Vines2018; Aoude, Haddad, Helset 2022;
Jakobsen, Mogull 2022; Jakobsen, Mogull, Plefka, Steinhoff 2021;



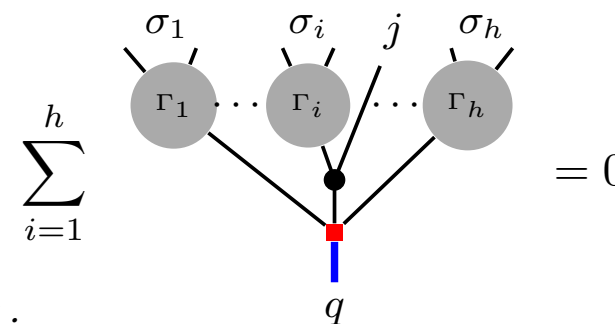
- KiHA for the interaction with fermions

e.g. Johansson, Ochirov 2016; GC, Johansson, Teng, Wang 2019

- Generalised CK and DC in the effective theory

e.g. Carrasco, Rodina, Yin, Zekioglu 2019,2021;
Chi, Elvang, Herderschee, Jones, Paranjape2022

$$\mathcal{L}_{\text{eff}} = \mathcal{L} + \sum_i c_i \mathcal{O}_i$$

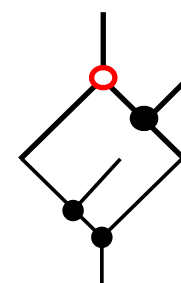


- KiHA in curved spacetime

e.g. Chacon, Nagy, White 2021; Monteiro, O'Connell, White2014 ; Herderschee, Roiban, Teng2022

- Loop level kinematic algebra

e.g. Borsten, Jurco, Kim, Macrelli, Saemann, Wolf 2020,2021,2022



“Thanks!”