



Jura, 1975

# About flavour, once again

Zoltan fest

ETH, Zurich, May 24, 2024

R. Barbieri  
SNS, Pisa



# Participating in the November Revolution

## Mixing of p Wave Axial Vector Resonances

#1

[Riccardo Barbieri \(CERN\)](#), [Raoul Gatto \(CERN\)](#), [Z. Kunszt \(Eotvos U.\)](#) (Dec, 1976)

Published in: *Phys.Lett.B* 66 (1977) 349-352



pdf



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cite



claim



reference search



11 citations

## Meson Masses and Widths in a Gauge Theory with Linear Binding Potential

#2

[Riccardo Barbieri \(CERN\)](#), [R. Kogerler \(CERN\)](#), [Z. Kunszt \(CERN\)](#), [Raoul Gatto \(Rome U.\)](#) (Jun, 1975)

Published in: *Nucl.Phys.B* 105 (1976) 125-138



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reference search



243 citations

## Electron-Positron Annihilation Above Charm Threshold

#3

[Riccardo Barbieri \(CERN\)](#), [R. Kogerler \(CERN\)](#), [Z. Kunszt \(CERN\)](#), [Raoul Gatto \(Rome U. and INFN, Rome\)](#) (Jun, 1975)

Published in: *Phys.Lett.B* 56 (1975) 477-481



pdf



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17 citations

## Meson hyperfine splittings and leptonic decays

#4

[Riccardo Barbieri \(CERN\)](#), [Raoul Gatto \(INFN, Rome\)](#), [R. Kogerler \(CERN\)](#), [Z. Kunszt \(CERN\)](#) (May, 1975)

Published in: *Phys.Lett.B* 57 (1975) 455-459



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reference search



264 citations

More to say at the end ...

# The SM Lagrangian (since 1973 in its full content)

$$\begin{aligned} \mathcal{L}_{\sim SM} = & -\frac{1}{4}F_{\mu\nu}^a F^{a\mu\nu} + i\bar{\Psi} \not{D}\Psi & (\sim 1975-2000) \\ & + |D_\mu h|^2 - V(h) & (\sim 1990-2012\text{-now}) \\ & + \psi_i \lambda_{ij} \psi_j h + h.c. & (\sim 2000\text{- now}) \end{aligned}$$

In () the approximate dates of the experimental confirmation of the various lines (at different levels)

The synthetic nature of the SM exhibited

0. Which rationale for matter quantum numbers?

$$\text{E.g.: } |Q_n - Q_p - Q_e| < 10^{-21} e$$

1. Phenomena unaccounted for

neutrino masses  
Dark matter

matter-antimatter asymmetry  
inflation?

2. Why  $\theta \lesssim 10^{-10}$  ?

$$\theta G_{\mu\nu} \tilde{G}^{\mu\nu}$$

Axions? A discrete space-time symmetry?

3.  $\mathcal{O}_i : d(\mathcal{O}_i) \leq 4$  only?

neutrino masses

Are the protons forever?

Gravity

What about individual  $L_i$  conservations?

4. Lack of calculability

⇒ the hierarchy problem  
the flavour puzzle

← none of the 15 masses  
predicted in the SM



# Where could some light come from?

1. A theory breakthrough

1a BSM

1b Foundations (FT, QM in curved space)

On 1a, not that one hasn't tried, sometimes with great ideas (GUT, susy, axion,...)

2. Astrophysics, Cosmology

2a DM, Dark Energy, B-asymmetry

2b Early Universe, Inflation

2c Black Holes, Grav. waves

Fundamental questions. Related to the structure of the SM of PP?

3. An experimental deviation  
from the SM

3a New particles

3b Precision

Focus on 3b, assuming (which requires) new physics in the MultiTeV,  
(as made likely by the hierarchy problem, still pending)

# An extreme summary in precision measurements

(European Strategy for PP, 2020)

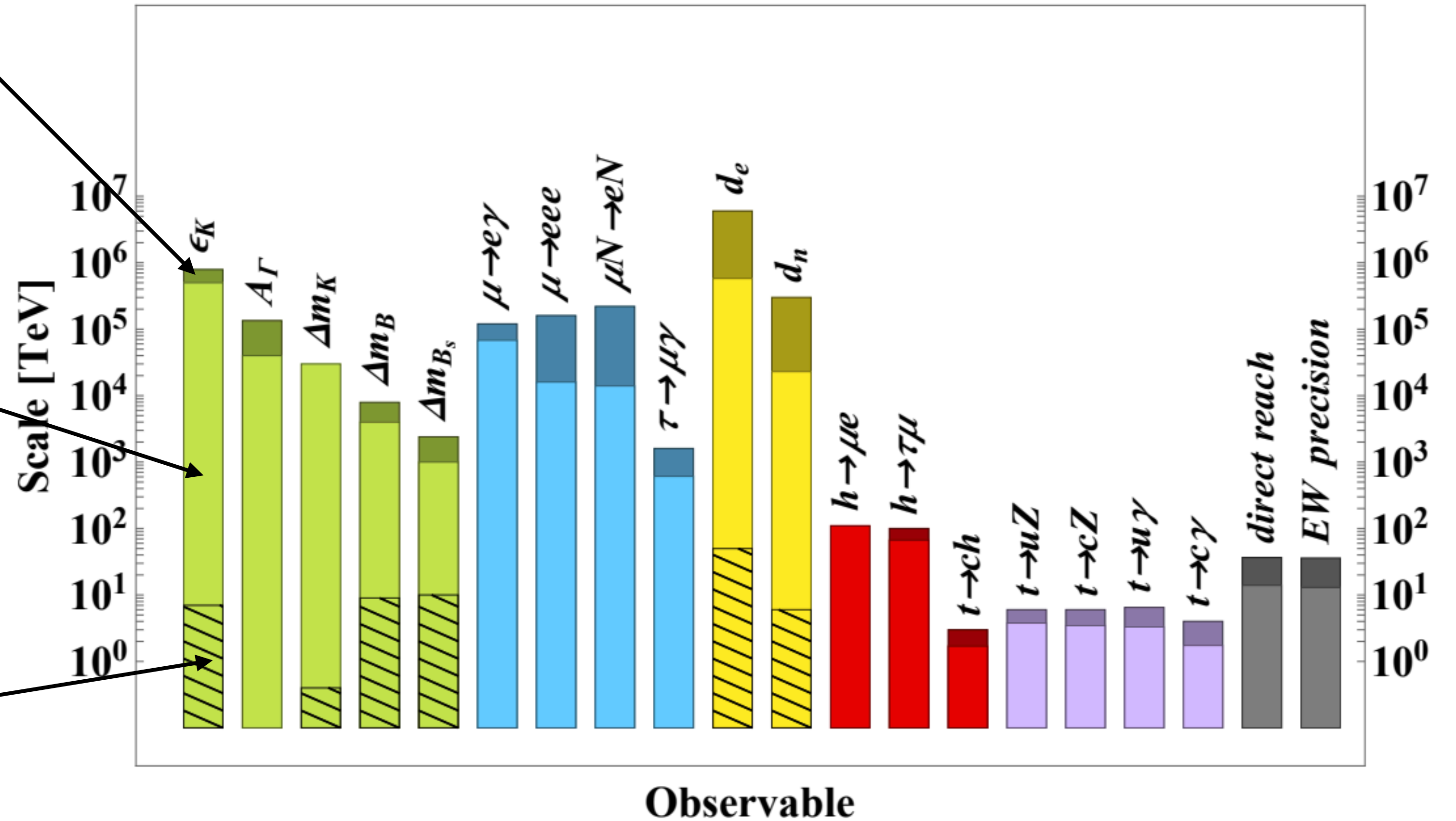
“Mid-term” prospects:  
All approved exp.s + LHCb II

current

$$\mathcal{L} = \mathcal{L}_{SM} + \sum_i \frac{1}{\Lambda_i^2} \mathcal{O}_i$$

current

$$\mathcal{L} = \mathcal{L}_{SM} + \sum_i \frac{F_i^{SM}}{\Lambda_i^2} \mathcal{O}_i$$



Flavour and EDMs (CPV) dominate



# A more detailed plot

(from actually measured flavour quantities only)

$$\mathcal{L} = \mathcal{L}_{SM} + \sum_i \frac{C_i}{\Lambda_i^2} \mathcal{O}_i$$

$$Q_1^{q_i q_j} = \bar{q}_{jL}^\alpha \gamma_\mu q_{iL}^\alpha \bar{q}_{jL}^\beta \gamma^\mu q_{iL}^\beta,$$

$$Q_2^{q_i q_j} = \bar{q}_{jR}^\alpha q_{iL}^\alpha \bar{q}_{jR}^\beta q_{iL}^\beta, \quad Q_4^{q_i q_j} = \bar{q}_{jR}^\alpha q_{iL}^\alpha \bar{q}_{jL}^\beta q_{iR}^\beta,$$

$$Q_3^{q_i q_j} = \bar{q}_{jR}^\alpha q_{iL}^\beta \bar{q}_{jR}^\beta q_{iL}^\alpha, \quad Q_5^{q_i q_j} = \bar{q}_{jR}^\alpha q_{iL}^\beta \bar{q}_{jL}^\beta q_{iR}^\alpha.$$

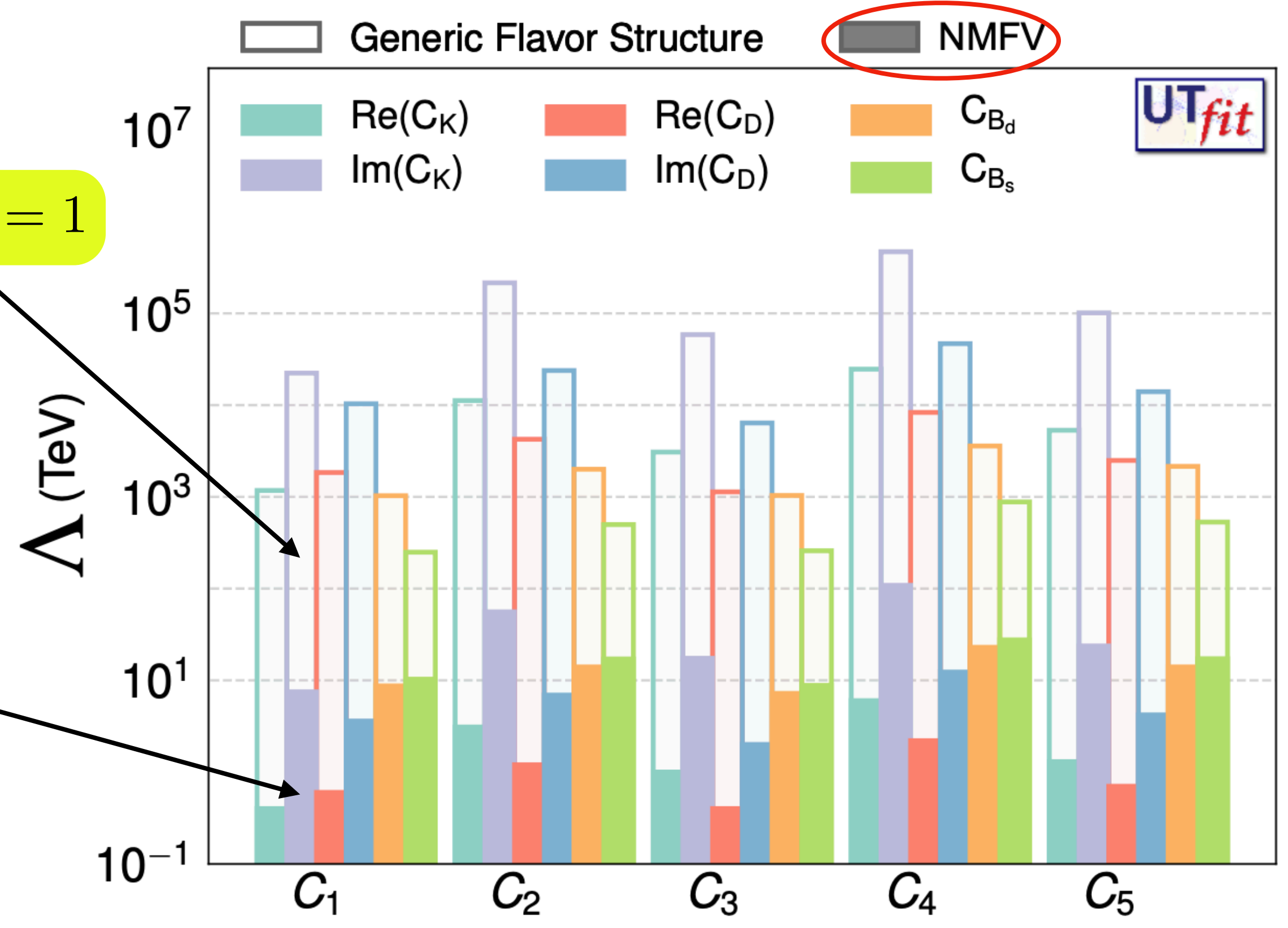
$$C_i = 1$$

NMFV

$$\mathcal{L} = \mathcal{L}_{SM} + \sum_i \frac{F_i^{SM}}{\Lambda_i^2} \mathcal{O}_i$$

$$F^{SM}(C_{K,D}) = (V_{td} V_{ts}^*)^2 e^{i\phi_{K,D}}$$

$$F^{SM}(C_{B_q}) = (V_{tq} V_{tb}^*)^2 e^{i\phi_{B_q}}$$



Pierini, 2023

Where is the actual scale of flavour physics  $\Lambda^f$ ? Must  $\Lambda^f$  be "low"?  
 How low can  $\Lambda^f$  be? Can one "make sense" of "NMFV"?

Must  $\Lambda^f$  be "low"?

Is the hierarchy problem related to the flavour puzzle?



# A difference in the two sectors of the SM?

$$\mathcal{L}_{SM} = -\frac{1}{4}F_{\mu\nu}^a F^{a\mu\nu} + i\bar{\Psi}\not{D}\Psi$$

The "gauge sector"

$$+|D_\mu\phi|^2 + M^2|\phi|^2 - \lambda|\phi|^4 + \Lambda + \lambda_{ij}\phi\bar{\Psi}_i\Psi_j$$

The "Higgs sector"

(where the Fermi scale originates)

the hierarchy  
problem

the CC problem

the flavour  
problem

In EFT they look  
much the same

No particle mass  
calculable (15=17-2)

To me: the relatively best motivation for BSM in the MultiTeV  
(and a strong motivation for the next HE collider)

# (Approximate) symmetries of the Yukawa couplings

Charged fermion Yukawa couplings

$$Y \propto U_L^+ \begin{pmatrix} m_1/m_3 & 0 & 0 \\ 0 & m_2/m_3 & 0 \\ 0 & 0 & 1 \end{pmatrix} U_R \quad m_1/m_3 \ll m_2/m_3 \ll 1 \quad U_L^u (U_L^d)^+ = V_{CKM} \approx \begin{pmatrix} 1 & \lambda & \lambda^3 \\ \lambda & 1 & \lambda^2 \\ \lambda^3 & \lambda^2 & 1 \end{pmatrix}$$

1 IF  $[U_L^{u,d}]_{i \neq j} \lesssim [V_{CKM}]_{i \neq j}$

$$Y^{u,d} \approx \begin{pmatrix} \text{light blue} & \text{light blue} & \text{light blue} \\ \text{medium blue} & \text{medium blue} & \text{medium blue} \\ \text{dark blue} & \text{dark blue} & \text{dark blue} \end{pmatrix} U(2)_q$$

$\Rightarrow U(2)_q$



# (approximate) symmetries of the Yukawa couplings

Charged fermion Yukawa couplings

$$Y \propto U_L^+ \begin{pmatrix} m_1/m_3 & 0 & 0 \\ 0 & m_2/m_3 & 0 \\ 0 & 0 & 1 \end{pmatrix} U_R \quad m_1/m_3 \ll m_2/m_3 \ll 1 \quad U_L^u (U_L^d)^+ = V_{CKM} \approx \begin{pmatrix} 1 & \lambda & \lambda^3 \\ \lambda & 1 & \lambda^2 \\ \lambda^3 & \lambda^2 & 1 \end{pmatrix}$$

1 IF  $[U_L^{u,d}]_{i \neq j} \lesssim [V_{CKM}]_{i \neq j}$

$$Y^{u,d} \approx \begin{pmatrix} \square & \square & \square \\ \square & \square & \square \\ \square & \square & \square \end{pmatrix} U(2)_q$$

$$\Rightarrow U(2)_q$$

2 IF  $1 + [U_R^{u,d}]_{i \neq j} \lesssim [U_L^{u,d}]_{i \neq j}$

$$Y^{u,d} \approx \begin{pmatrix} \square & \square & \square \\ \square & \square & \square \\ \square & \square & \square \end{pmatrix} U(2)_q$$

$U(2)_{u,d}$

$$\Rightarrow U(2)_q \times U(2)_u \times U(2)_d$$

# (approximate) symmetries of the Yukawa couplings

Charged fermion Yukawa couplings

$$Y \propto U_L^+ \begin{pmatrix} m_1/m_3 & 0 & 0 \\ 0 & m_2/m_3 & 0 \\ 0 & 0 & 1 \end{pmatrix} U_R \quad m_1/m_3 \ll m_2/m_3 \ll 1 \quad U_L^u (U_L^d)^+ = V_{CKM} \approx \begin{pmatrix} 1 & \lambda & \lambda^3 \\ \lambda & 1 & \lambda^2 \\ \lambda^3 & \lambda^2 & 1 \end{pmatrix}$$

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$$Y^{u,d} \approx \begin{pmatrix} \square & \square & \square \\ \square & \square & \square \\ \square & \square & \square \end{pmatrix} U(2)_q$$

$U(2)_{u,d}$

$$\Rightarrow U(2)_q \times U(2)_u \times U(2)_d$$

3 IF  $2 + [U_{L,R}^e]_{i \neq j} \lesssim [U_{L,R}^{u,d}]_{i \neq j}$

$$\Rightarrow U(2)_q \times U(2)_u \times U(2)_d \times U(2)_l \times U(2)_e$$

Can  $U(2)^n$  emerge as an accidental symmetry?  
What breaks it?

B, Isidori et al, 2011

# A definite goal: Precision in composite Higgs

What is the radius of Higgs compositeness, if any?  $l_H = 1/m_*$

A two-parameter  
"theory"

_____	$m_* = g_* f$
_____	$f$
_____	$m_H$

Giudice et al, 2007

$H = \text{pNGB}$

$f$  = scale of symmetry breaking  
 $m_*$  = scale of Higgs compositeness

Fine tuning =  $(\frac{v}{f})^2$   $v = 175\text{GeV}$



# An EFT approach

$$\mathcal{L} = \mathcal{L}_{SM} + \frac{m_*^4}{g_*^2} \mathcal{L}_{res} \left( \frac{g_* \Phi^d}{m_*^d}, \frac{d_\mu}{m_*}, \frac{g A_\mu}{m_*}, \frac{\lambda_\Psi^i \Psi^i}{m_*^{3/2}} \right)$$

Giudice et al, 2007

$\Phi^d$  = Strong resonances of dim d (J=0,1/2,1,...) including  $H$

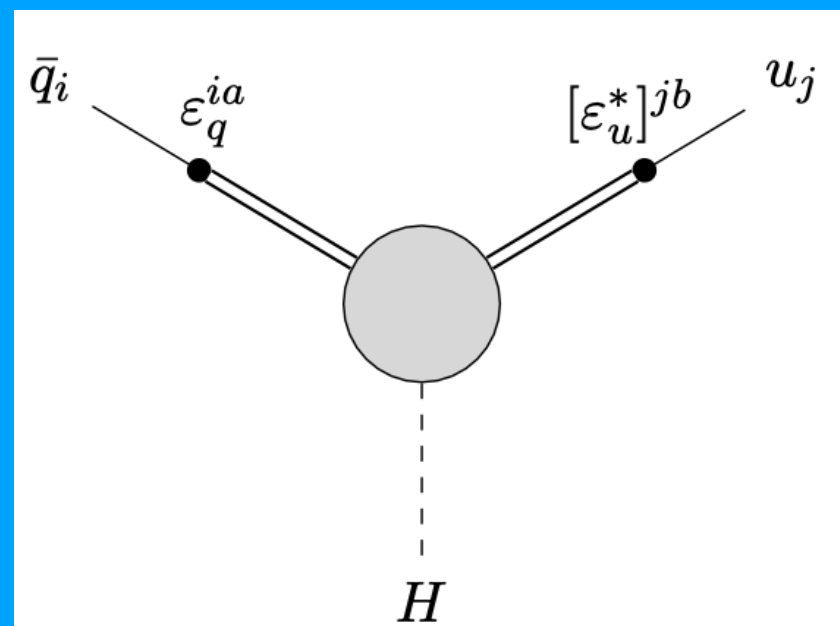
$\Psi^i, A_\mu$  = SM fields

$\lambda_\Psi^i$  = flavour par.s, subject to suitable symmetries

Redi, Wyler 2011

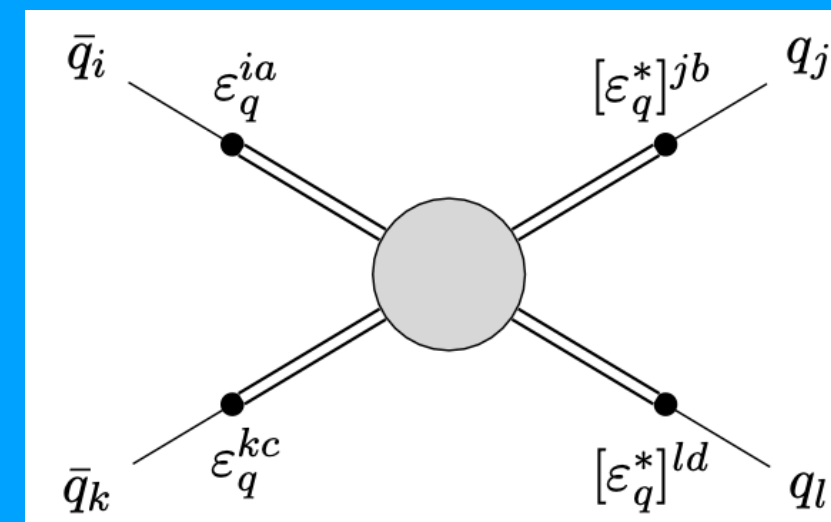
B, Buttazzo et al, 2013

Glioti et al, 2024



$$Y_{ij}^u = g_* \epsilon_q^{ia} C_{ab} [\epsilon_u^*]^{jb}$$

E.g.



$$\mathcal{L}^{4q} = \frac{g_*^2}{m_*^2} C_{abcd} \epsilon_q^{ia} [\epsilon_q^*]^{jb} \epsilon_q^{kc} [\epsilon_q^*]^{ld} \bar{q}^i \gamma_\mu q^j \bar{q}^k \gamma^\mu q^l$$

$$\epsilon_f^{ia} = \frac{\lambda_f^{ia}}{g_*}$$

$$C_{ab}, C_{abcd} = \mathcal{O}(1)$$

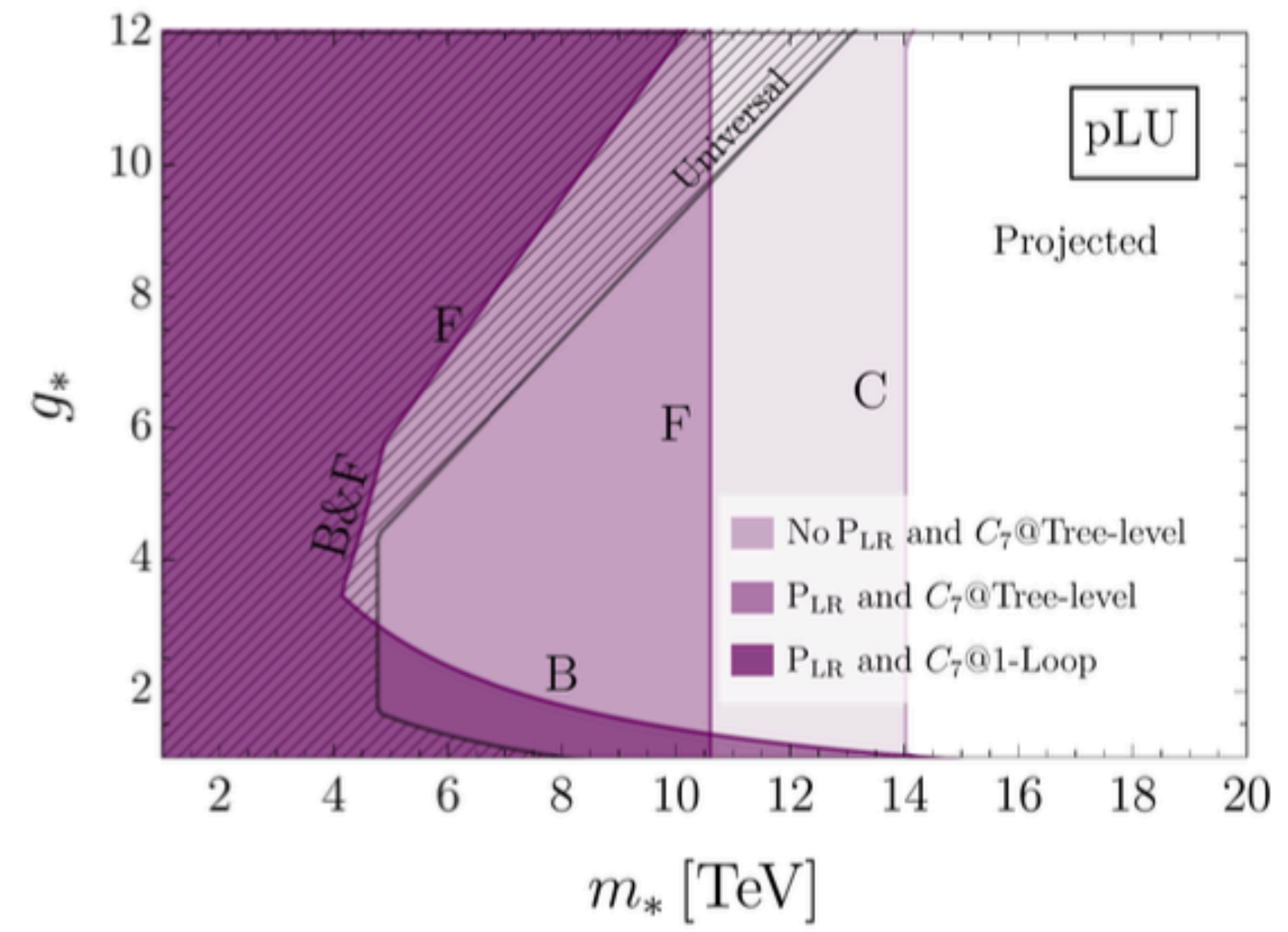
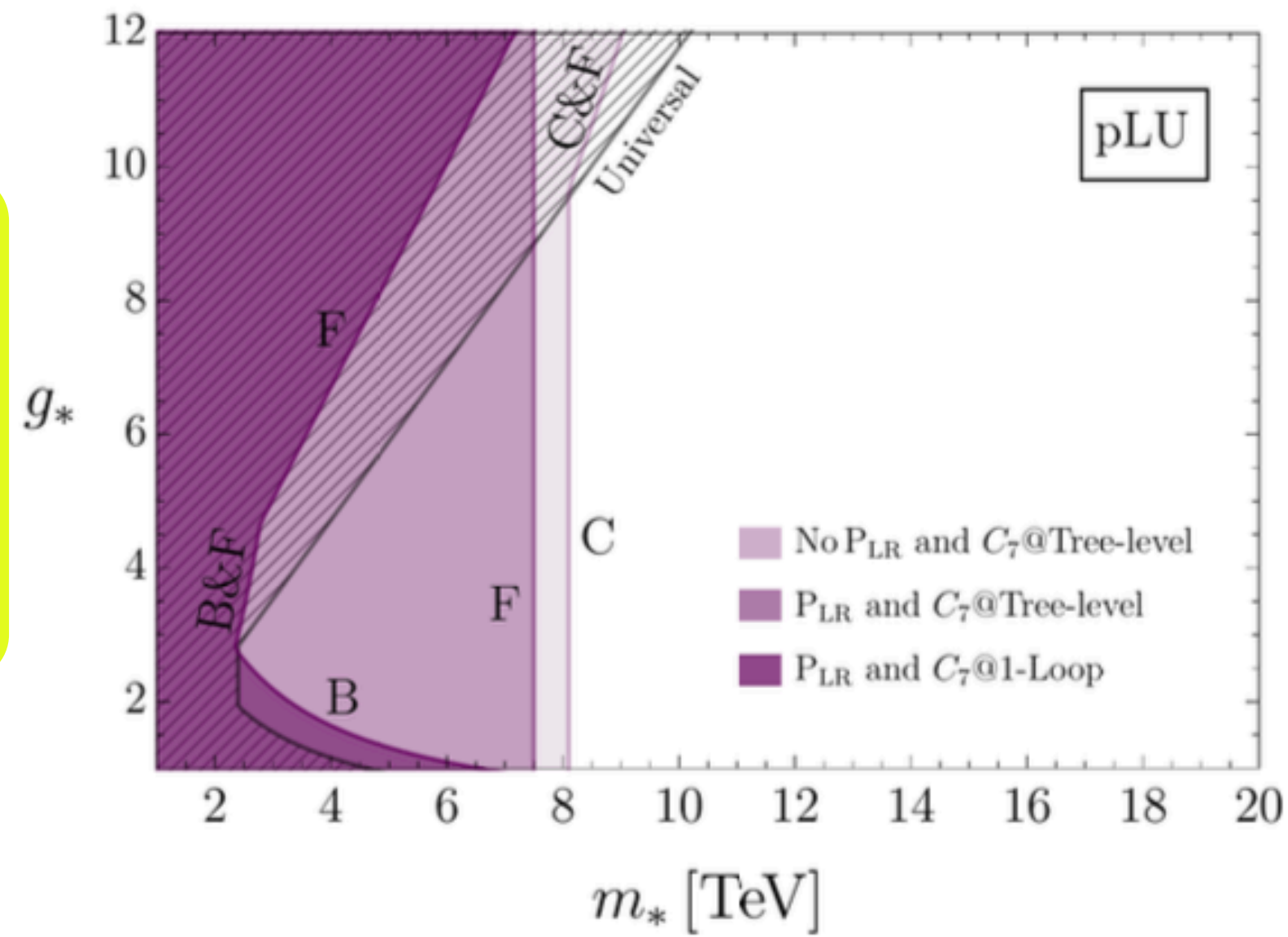
If  $\epsilon_{11} \ll \epsilon_{22} \ll \epsilon_{33}$  "anarchy"  $\Rightarrow m_* \gtrsim 10^{2 \div 3} TeV (g_*/4\pi)$



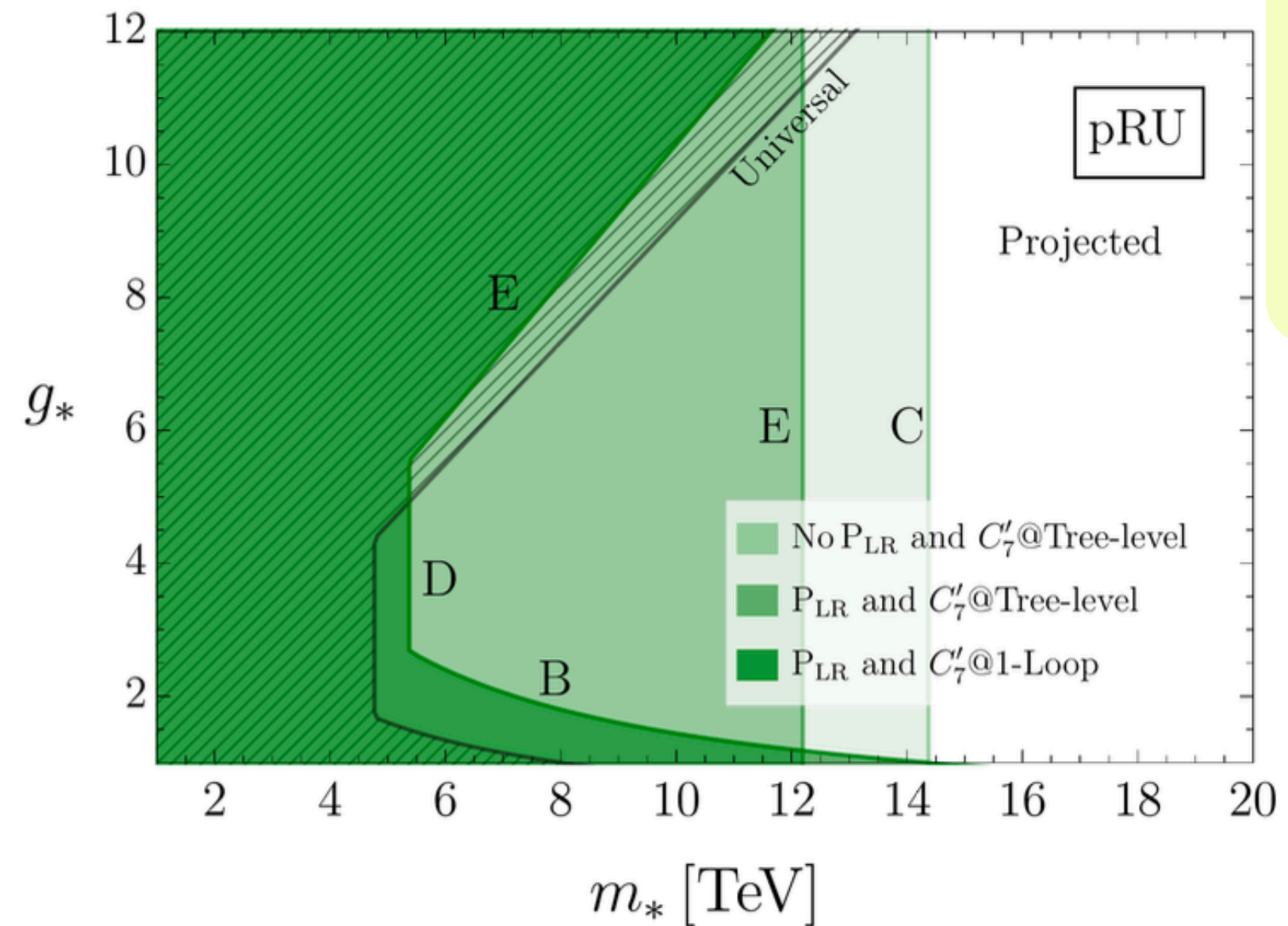
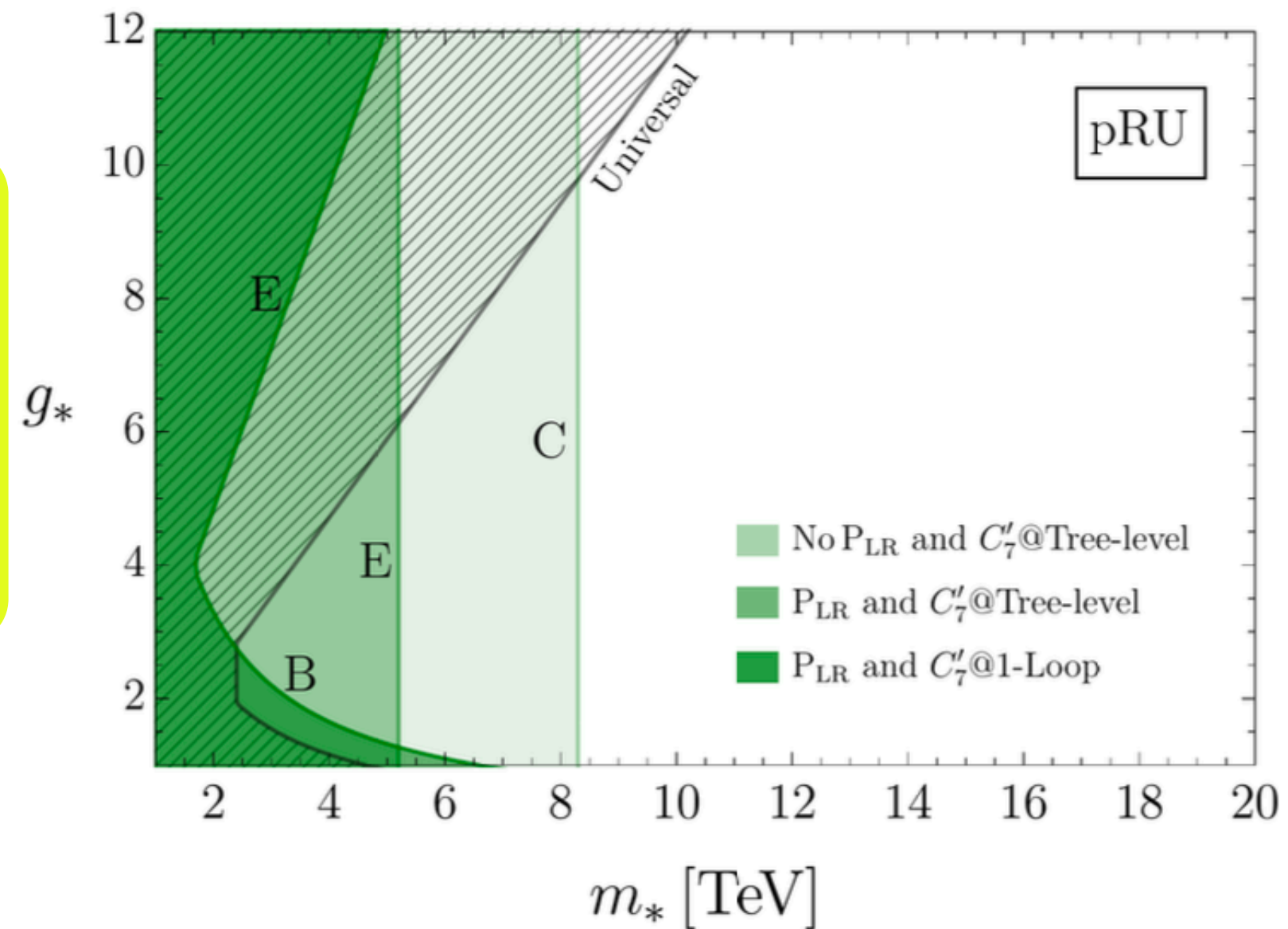
# Summary of excluded/sensitivity regions

Glioti et al, 2024

partial Left Univ =  
 $\lambda_{\Psi}^i$  respecting  
 $U(2)_q$



part. Right Univ =  
 $\lambda_{\Psi}^i$  respecting  
 $U(2)_u \times U(2)_d$



Label	Observable
A	$pp \rightarrow jj$
B	$\Delta F = 2 (B_d)$
C	$B_s \rightarrow \mu^+ \mu^-$
D	nEDM
E	$B^0 \rightarrow K^{*0} e^+ e^- (C'_7)$
F	$B \rightarrow X_s \gamma (C_7)$
G	W-coupling

Universal = flavour-less EW observables

Projected = "mid-term"

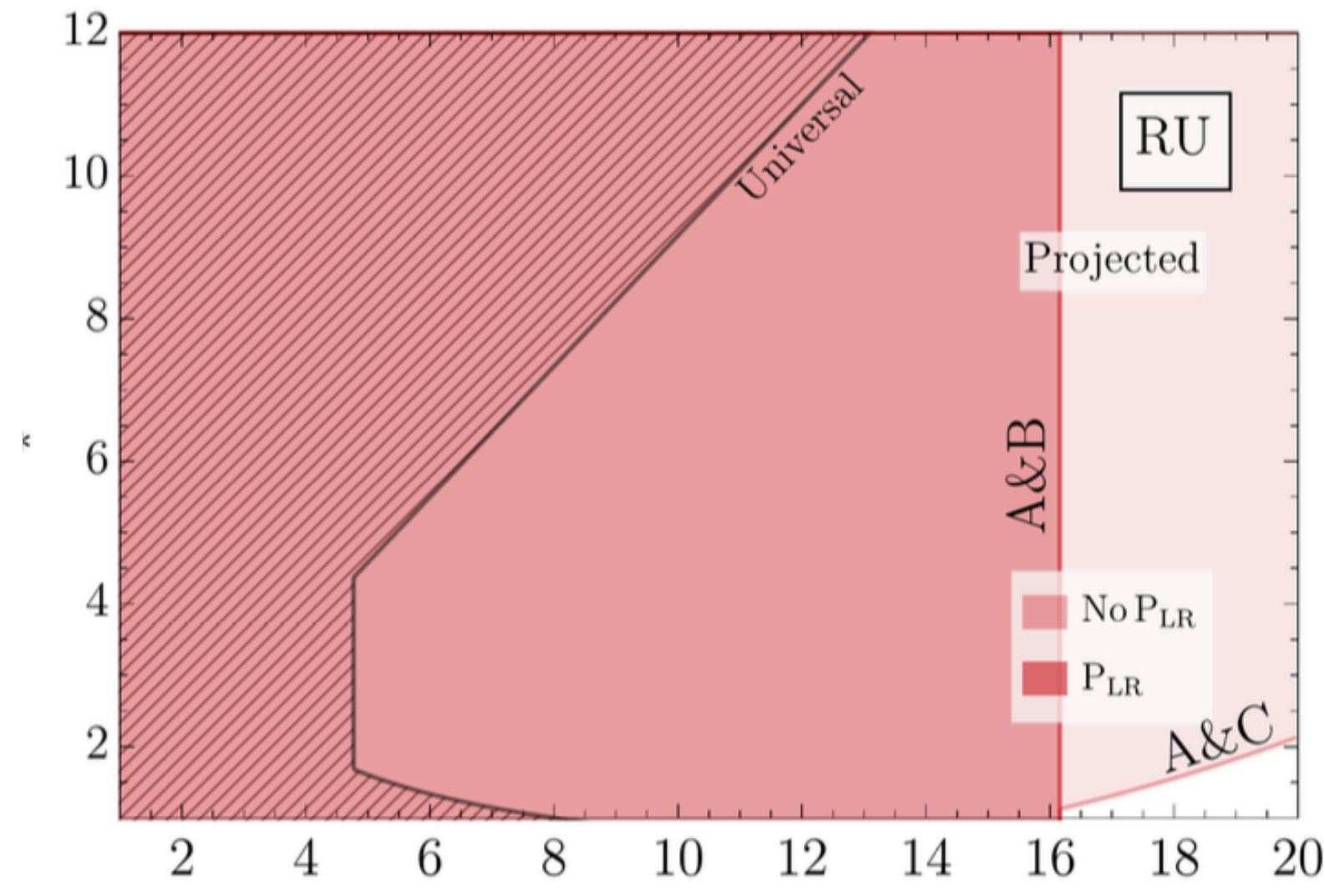
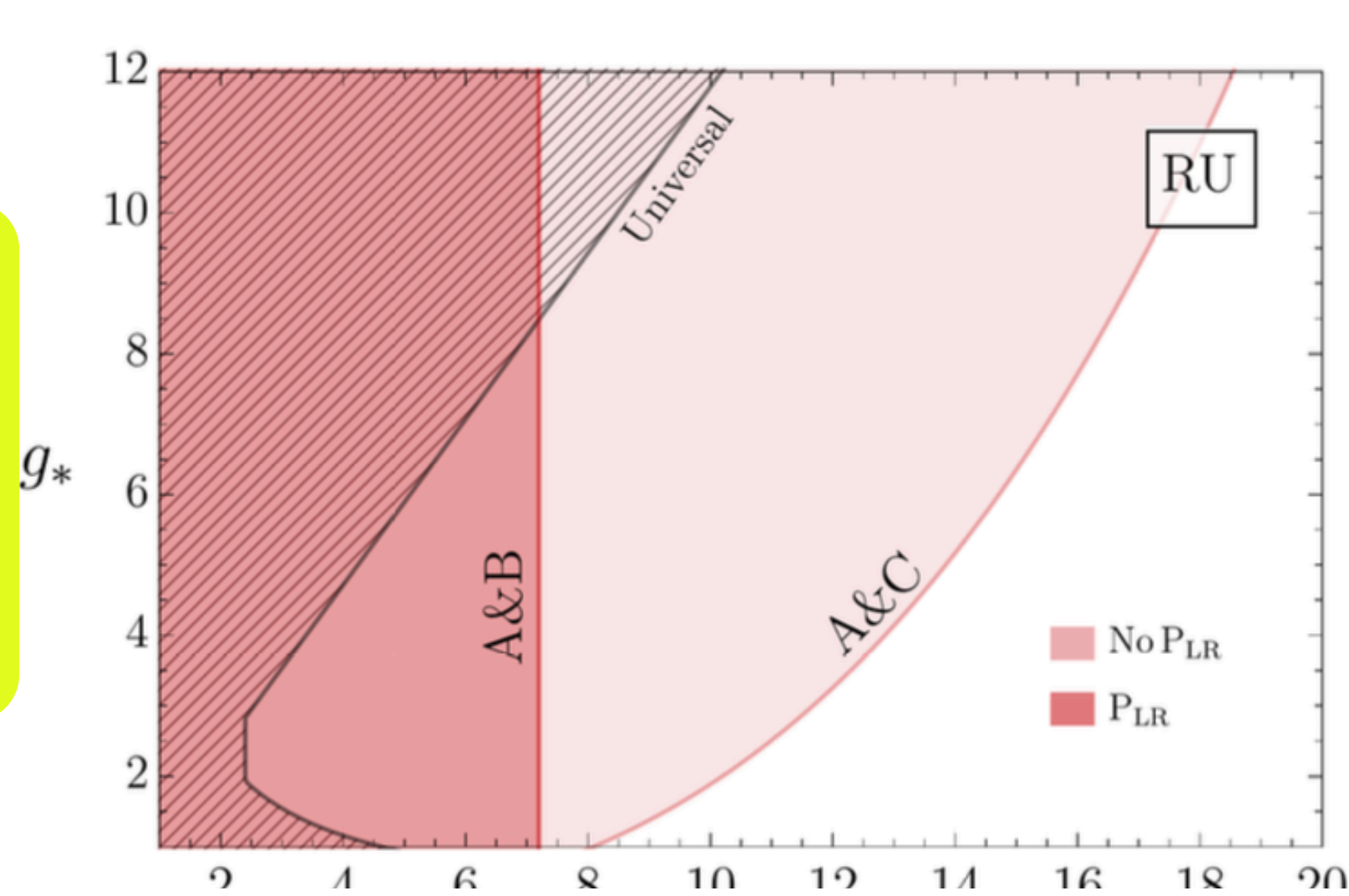
(Quarks only)



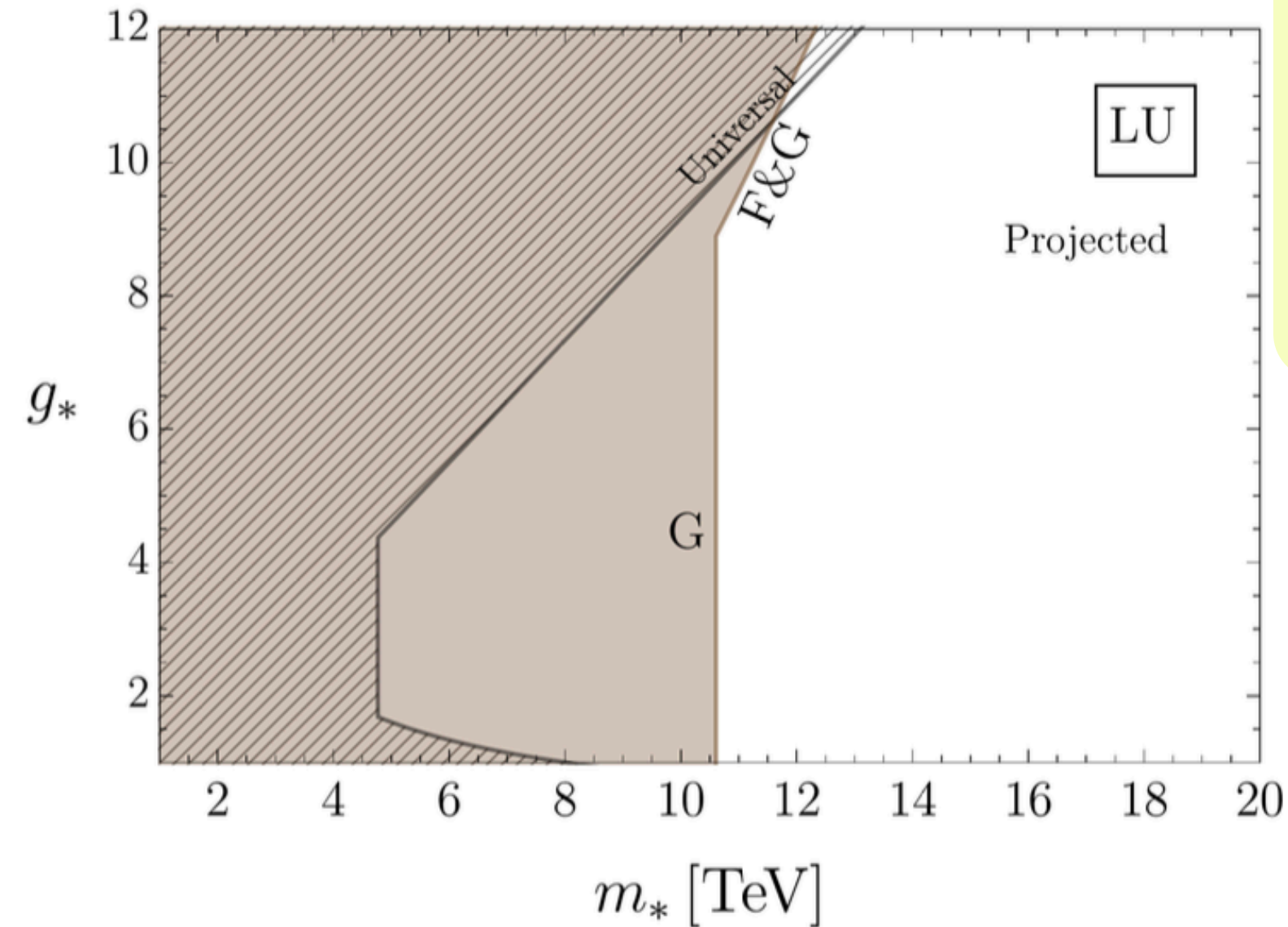
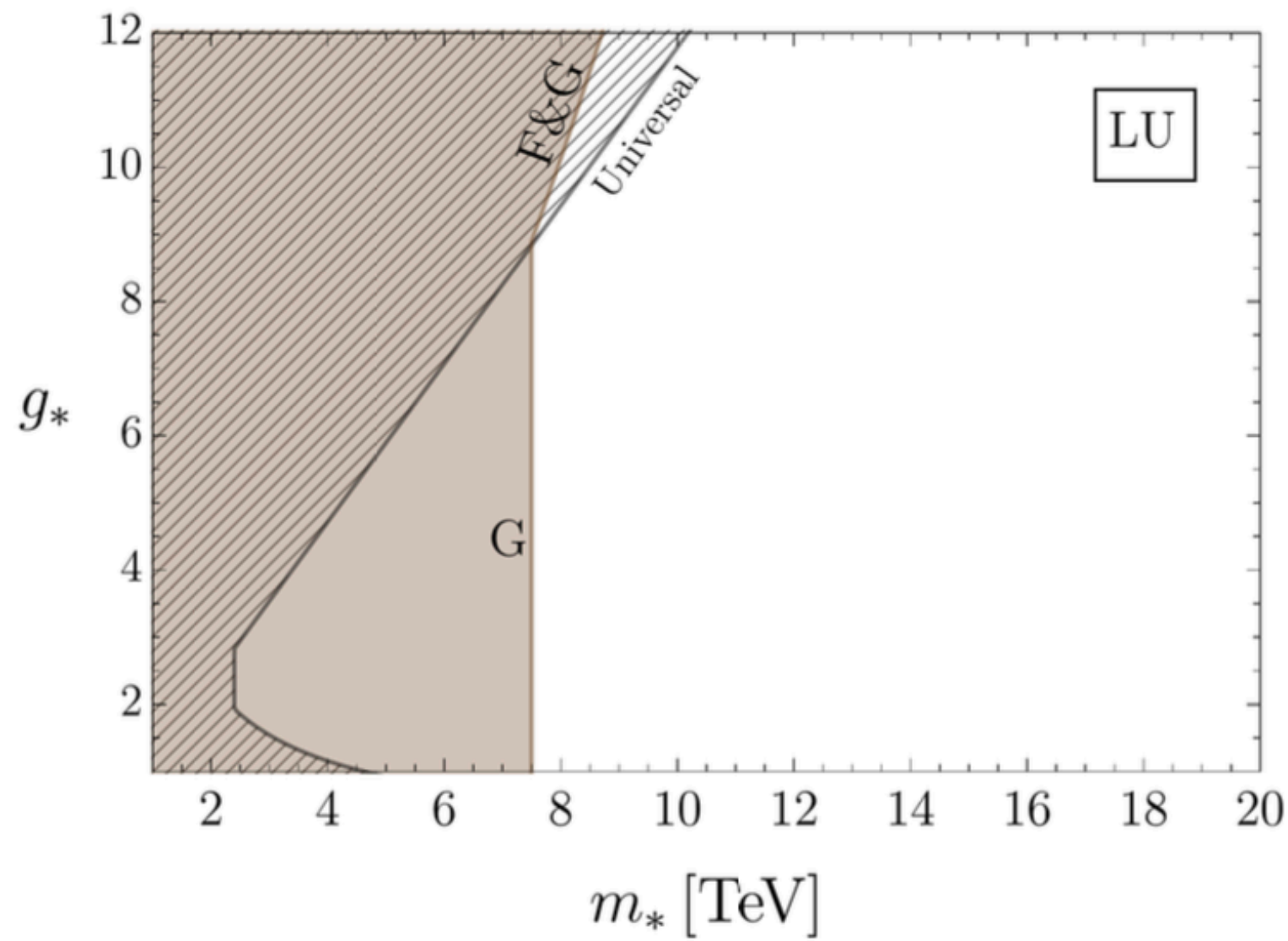
# Summary of excluded/sensitivity regions

Glioti et al, 2024

Right Univ =  
 $\lambda_{\Psi}^i$  respecting  
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Left Univ =  
 $\lambda_{\Psi}^i$  respecting  
 $U(3)_q$



Label	Observable
A	$pp \rightarrow jj$
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F	$B \rightarrow X_s \gamma (C_7)$
G	W-coupling

Mid-term prospect:  $m_* > (11 \div 20) TeV$  for any  $g_*$

(Quarks only)



# Can one "make sense" of "NMFV"?

(or of the symmetries respected by  $\lambda_{\Psi}^i$  in the previous case?)

$$\mathcal{L} = \mathcal{L}_{SM} + \sum_i \frac{F_i^{SM}}{\Lambda_i^2} \mathcal{O}_i$$

$$F^{SM}(C_{K,D}) = (V_{td}V_{ts}^*)^2 e^{i\phi_{K,D}}$$

$$F^{SM}(C_{B_q}) = (V_{tq}V_{tb}^*)^2 e^{i\phi_{B_q}}$$

Can one relate the  $F_i^{SM}$  to the structure of  $Y^{u,d}$ ?

# All the SM in 1 page

## 1. Symmetry group $L \times \mathcal{G}$

$L$  = Lorentz (space-time)

$\mathcal{G} = SU(3) \times SU(2) \times U(1)$  (local)

## 2. Particle content (rep.s of $L \times \mathcal{G}$ )

	$h$	$Q$	$L$	$u$	$d$	$e$
Lorentz	0	$1/2_L$	$1/2_L$	$1/2_R$	$1/2_R$	$1/2_R$
$SU(3)$	<b>1</b>	<b>3</b>	<b>1</b>	<b>3</b>	<b>3</b>	<b>1</b>
$SU(2)$	<b>2</b>	<b>2</b>	<b>2</b>	<b>1</b>	<b>1</b>	<b>1</b>
$U(1)$	$-1/2$	$1/6$	$-1/2$	$2/3$	$-1/3$	$-1$

## 3. All "operators" (products of $\Phi, \partial_\mu \Phi$ ) in $\mathcal{L}$ of dimension $\leq 4$

$$\hbar = c = 1 \Rightarrow [A_\mu] = [\phi] = [\partial_\mu] = M, \quad [\Psi] = M^{3/2}, \quad [\mathcal{L}] = M^4$$

# Minimal Flavour Deconstruction

B, Isidori, 2023

$$SU(3) \times SU(2) \times G_Y$$

$$G_Y = U(1)_Y^{[3]} \times U(1)_{B-L}^{[12]} \times U(1)_{T_{3R}}^{[2]} \times U(1)_{T_{3R}}^{[1]} \quad H \stackrel{G_Y}{=} (-1/2, 0, 0, 0)$$

$$G_Y \xrightarrow{\sigma} U(1)_Y^{[3]} \times U(1)_{B-L}^{[12]} \times U(1)_{T_{3R}}^{[12]} \xrightarrow{\phi, \chi} U(1)_Y$$

$$\epsilon_\sigma = \frac{\langle \sigma \rangle}{\Lambda_{[12]}}$$

$$\epsilon_\phi = \frac{\langle \phi \rangle}{\Lambda_{[23]}}$$

$$\epsilon_\chi = \frac{\langle \chi^{q,l} \rangle}{\Lambda_{[23]}}$$

$$Y \sim \left( \begin{array}{c|c|c} & U(1)_{B-L}^{[12]} & \\ \hline & U(1)_{T_{3R}}^{[1]} & U(1)_{T_{3R}}^{[2]} \\ \hline O(\epsilon_\sigma \epsilon_\phi) & O(\epsilon_\phi) & O(\epsilon_\chi) \\ \hline O(\epsilon_\sigma \epsilon_\phi \epsilon_\chi) & O(\epsilon_\phi \epsilon_\chi) & O(1) \end{array} \right) U(1)_{B-L}^{[12]}$$

(Still EFT)

Can one construct an explicit 4d gauge theory without small Yukawa couplings?

(Where do the  $\Lambda$ 's come from?)



# Minimal Flavour Deconstruction in 4d

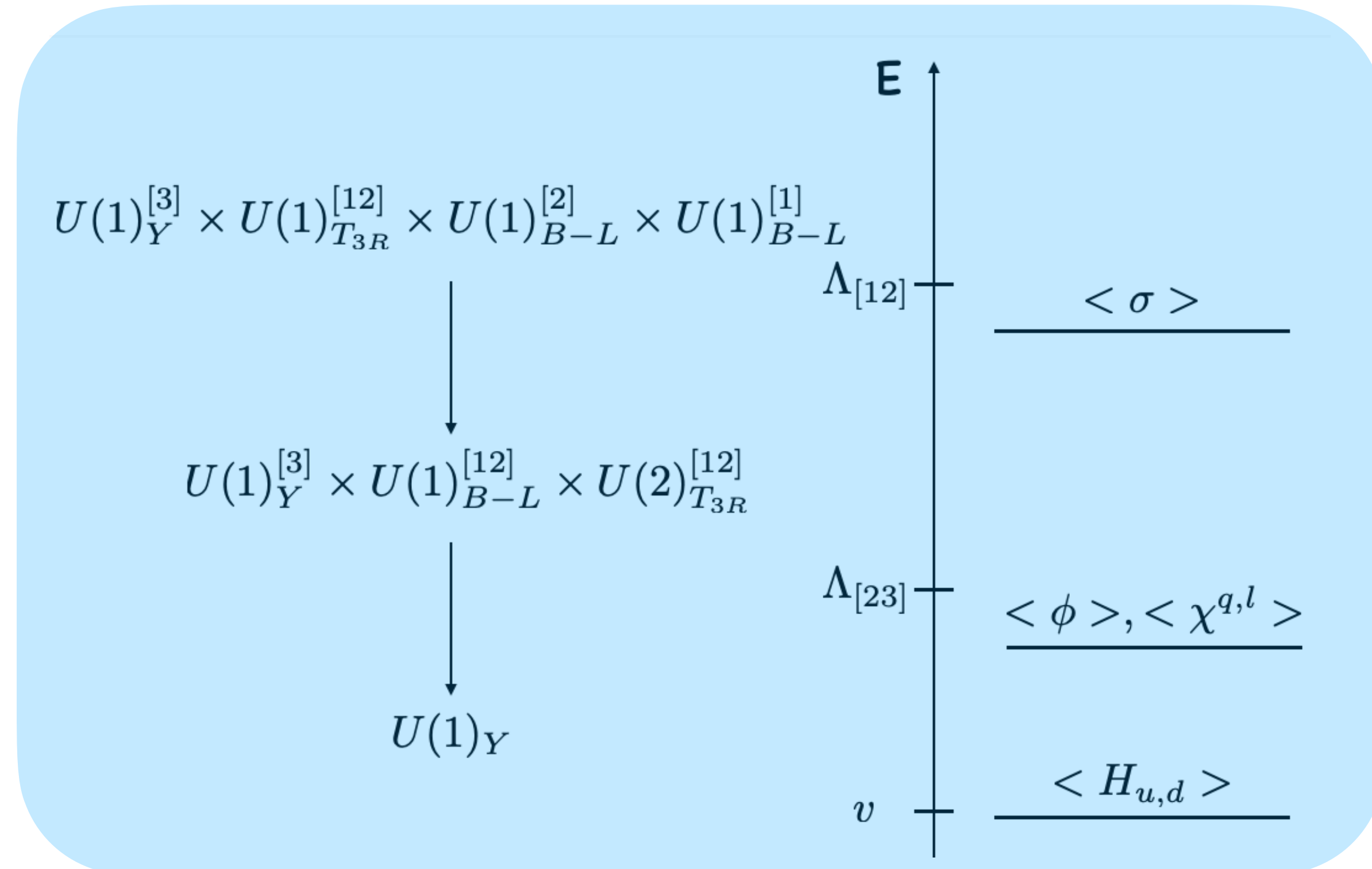
vev scale	Field	$U(1)_Y^{[3]}$	$U(1)_{B-L}^{[12]}$	$U(1)_{T_{3R}}^{[2]}$	$U(1)_{T_{3R}}^{[1]}$	$SU(3) \times SU(2)$
$v$	$H_{u,d}$	$-1/2$	$0$	$0$	$0$	$(\mathbf{1}, \mathbf{2})$
$O(10^{-1}) \times \Lambda_{[23]}$	$\chi^q$	$-1/6$	$1/3$	$0$	$0$	$(\mathbf{1}, \mathbf{1})$
	$\chi^l$	$1/2$	$-1$	$0$	$0$	$(\mathbf{1}, \mathbf{1})$
	$\phi$	$1/2$	$0$	$-1/2$	$0$	$(\mathbf{1}, \mathbf{1})$
$O(10^{-1}) \times \Lambda_{[12]}$	$\sigma$	$0$	$0$	$1/2$	$-1/2$	$(\mathbf{1}, \mathbf{1})$

$$Z_2 : H_u \rightarrow u, H_d \rightarrow d, e$$

$$\tan\beta = v_u/v_d = 10 \div 30$$

$$V = \lambda(\chi^q)^3 \chi^l$$

Universal breaking  
of the gauge group



$$\epsilon_\sigma = \frac{\langle \sigma \rangle}{\Lambda_{[12]}}$$

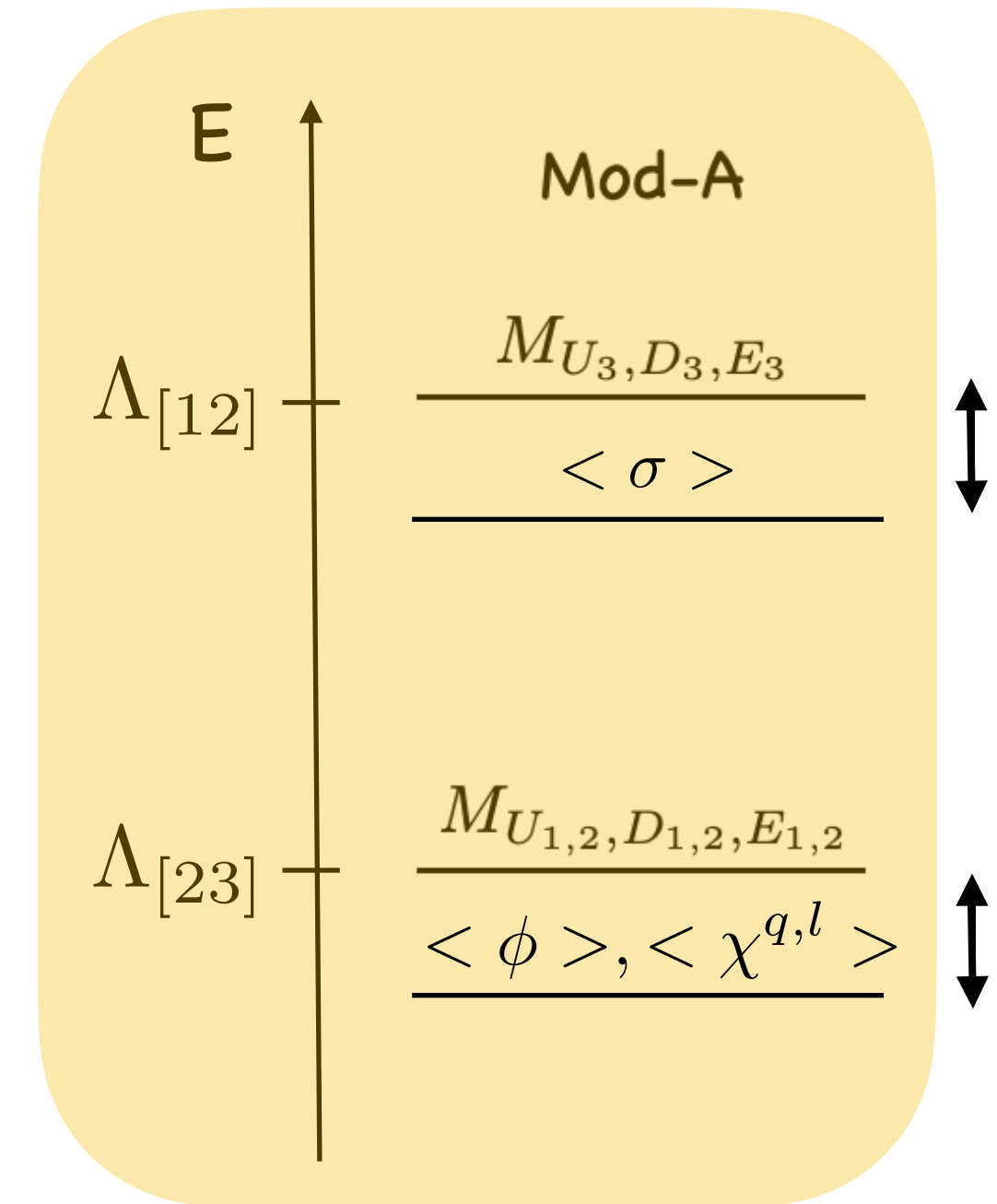
$$\epsilon_\phi = \frac{\langle \phi \rangle}{\Lambda_{[23]}}, \quad \epsilon_\chi = \frac{\langle \chi^{q,l} \rangle}{\Lambda_{[23]}}$$

(Where do the  $\Lambda$ 's come from?)

# Model A 1st way for the $\Lambda$ 's

Vector-like fermions

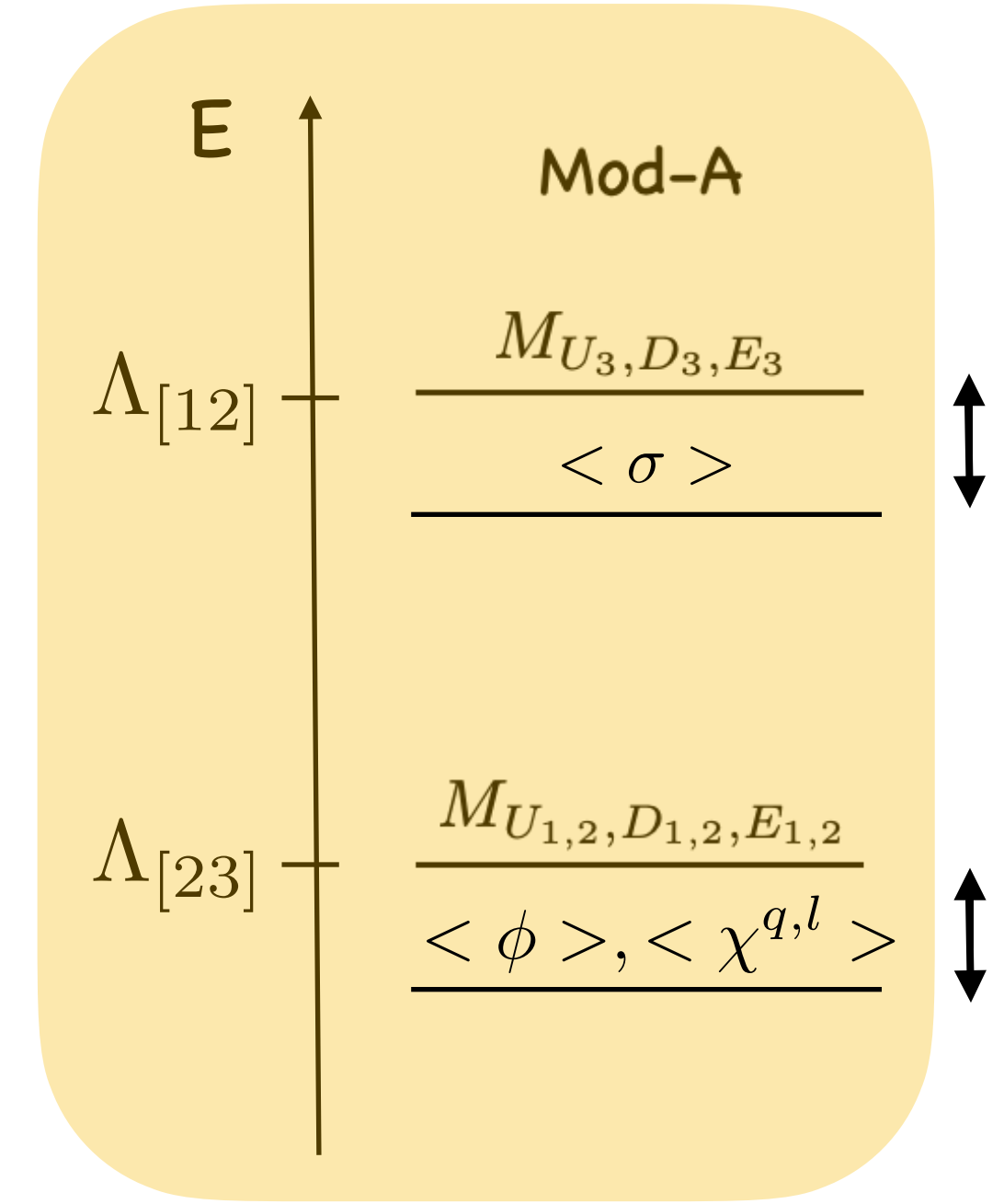
		$U(1)_Y^{[3]}$	$U(1)_{B-L}^{[12]}$	$U(1)_{T_{3R}}^{[2]}$	$U(1)_{T_{3R}}^{[1]}$	$SU(3) \times SU(2)$
light VL ( $\alpha = 1, 2$ )	$U_\alpha$	1/2	1/3	0	0	$(\mathbf{3}, \mathbf{1})$
	$D_\alpha$	-1/2	1/3	0	0	$(\mathbf{3}, \mathbf{1})$
	$E_\alpha$	-1/2	-1	0	0	$(\mathbf{1}, \mathbf{1})$
heavy VL	$U_3$	0	1/3	1/2	0	$(\mathbf{3}, \mathbf{1})$
	$D_3$	0	1/3	-1/2	0	$(\mathbf{3}, \mathbf{1})$
	$E_3$	0	-1	-1/2	0	$(\mathbf{1}, \mathbf{1})$



# Model A 1st way for the $\Lambda$ 's

Vector-like fermions

		$U(1)_Y^{[3]}$	$U(1)_{B-L}^{[12]}$	$U(1)_{T_{3R}}^{[2]}$	$U(1)_{T_{3R}}^{[1]}$	$SU(3) \times SU(2)$
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heavy VL	$U_3$	0	1/3	1/2	0	$(\mathbf{3}, \mathbf{1})$
	$D_3$	0	1/3	-1/2	0	$(\mathbf{3}, \mathbf{1})$
	$E_3$	0	-1	-1/2	0	$(\mathbf{1}, \mathbf{1})$



Most general  $d \leq 4$

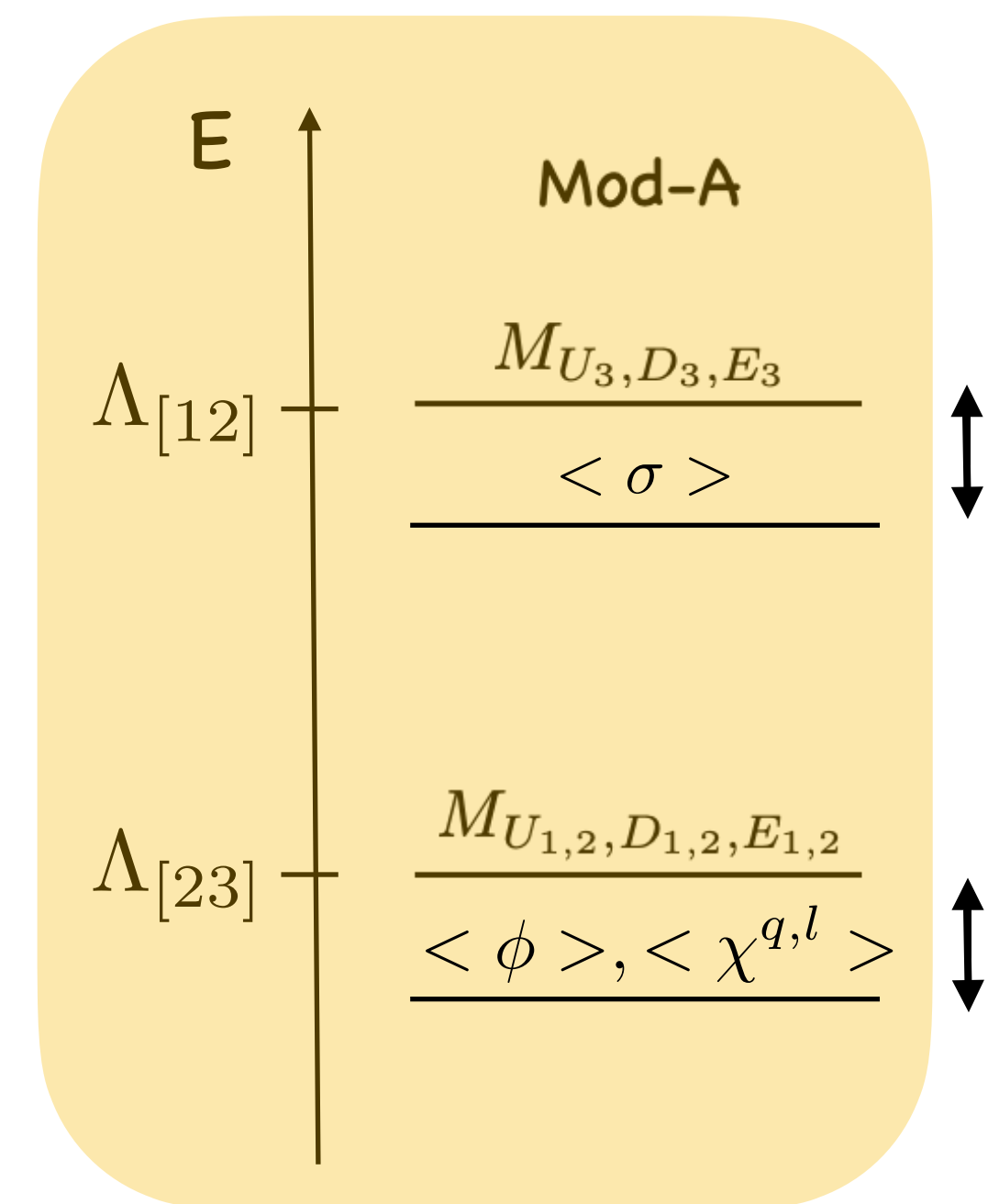
$$\mathcal{L}_Y^u = \left( \underbrace{y_3^u \bar{q}_3 u_3 H_u}_{\text{circled}} + y_{i\alpha}^u \bar{q}_i U_\alpha H_u + y_\alpha^{\chi^u} \bar{U}_\alpha u_3 \chi^q + y_{\alpha 2}^{\phi_u} \bar{U}_\alpha u_2 \phi + y_{\alpha 3}^{\phi_u} \bar{U}_{R\alpha} U_{L3} \phi + \hat{y}_{\alpha 3}^{\phi_u} \bar{U}_{L\alpha} U_{R3} \phi + y_1^{\sigma u} \bar{U}_3 u_1 \sigma + \text{h.c.} \right) + M_{U_3} \bar{U}_3 U_3 + M_{U_\alpha} \bar{U}_\alpha U_\alpha$$



# Model A 1st way for the $\Lambda$ 's

Vector-like fermions

		$U(1)_Y^{[3]}$	$U(1)_{B-L}^{[12]}$	$U(1)_{T_{3R}}^{[2]}$	$U(1)_{T_{3R}}^{[1]}$	$SU(3) \times SU(2)$
light VL ( $\alpha = 1, 2$ )	$U_\alpha$	1/2	1/3	0	0	( <b>3</b> , 1)
	$D_\alpha$	-1/2	1/3	0	0	( <b>3</b> , 1)
	$E_\alpha$	-1/2	-1	0	0	( <b>1</b> , 1)
heavy VL	$U_3$	0	1/3	1/2	0	( <b>3</b> , 1)
	$D_3$	0	1/3	-1/2	0	( <b>3</b> , 1)
	$E_3$	0	-1	-1/2	0	( <b>1</b> , 1)

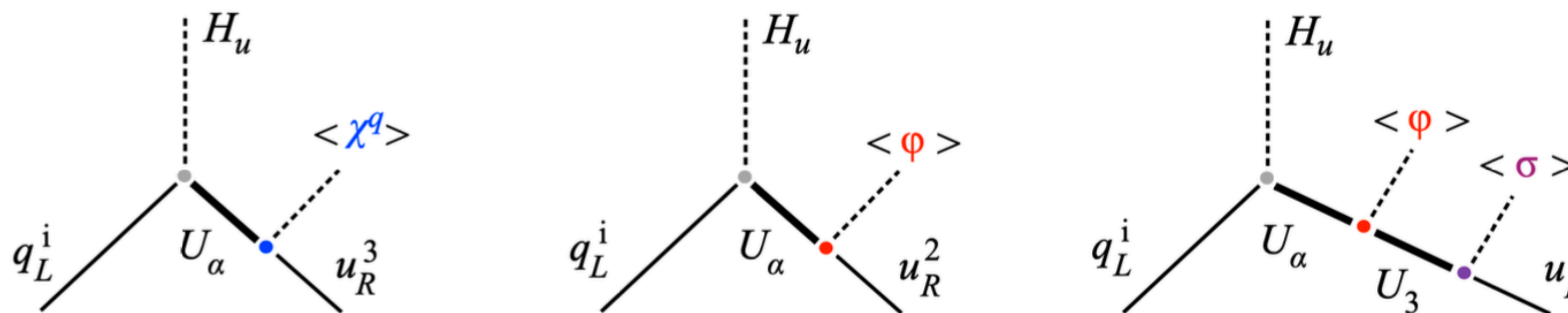


Most general  $d \leq 4$

$$\mathcal{L}_Y^u = \left( y_3^u \bar{q}_3 u_3 H_u + y_{i\alpha}^u \bar{q}_i U_\alpha H_u + y_\alpha^{\chi^u} \bar{U}_\alpha u_3 \chi^q + y_{\alpha 2}^{\phi_u} \bar{U}_\alpha u_2 \phi + y_{\alpha 3}^{\phi_u} \bar{U}_{R\alpha} U_{L3} \phi + \hat{y}_{\alpha 3}^{\phi_u} \bar{U}_{L\alpha} U_{R3} \phi + y_1^{\sigma_u} \bar{U}_3 u_1 \sigma + \text{h.c.} \right) + M_{U_3} \bar{U}_3 U_3 + M_{U_\alpha} \bar{U}_\alpha U_\alpha$$

$Y_u$

$\Rightarrow$



And similarly  
for  $Y_{d,e}$

$$Y_u \approx \begin{pmatrix} y_{1\alpha}^u \hat{y}_{\alpha 3}^{\phi_u} y_1^{\sigma_u} \epsilon_\sigma \epsilon_\phi & y_{1\alpha}^u y_{\alpha 2}^{\phi_u} \epsilon_\phi & y_{12}^u y_2^{\chi^u} \epsilon_\chi \\ y_{2\alpha}^u \hat{y}_{\alpha 3}^{\phi_u} y_1^{\sigma_u} \epsilon_\sigma \epsilon_\phi & y_{2\alpha}^u y_{\alpha 2}^{\phi_u} \epsilon_\phi & y_{22}^u y_2^{\chi^u} \epsilon_\chi \\ \approx 0 & \approx 0 & y_3^u \end{pmatrix}$$

$$\frac{v_2}{v_1} \approx 10 \quad \epsilon_\chi \approx \epsilon_\phi \approx 5 \cdot 10^{-2}$$

$$y'_s = 0.1 \div 1 \quad \epsilon_\sigma \approx 2 \cdot 10^{-2}$$

# Phenomenology at $\Lambda_{[23]}$ (universal)

$$U(1)_Y^{[3]} \times U(1)_{(B-L)/2}^{[12]} \times U(1)_{T_{3R}}^{[2]} \times U(1)_{T_{3R}}^{[1]}$$

$$g_3 \qquad g_B \qquad g_T \qquad g_T$$

$$g_3 \gg g_B, g_T$$

$$g_B = g' / \sin\alpha$$

$$g_T = g' / \cos\alpha$$

For every tree level effect at  $\Lambda_{[23]}$ , 4 parameters:  $m_{Z_{23}}, \alpha, b$  and  $(\theta_L^d)_{23}$

$$U(1)_Y^{[3]} \times U(1)_{B-L}^{[12]} \times U(1)_{T_{3R}}^{[12]}$$

$$\Lambda_{[23]} \quad \frac{Z'_{23}}{Z_{23}}$$

$$Z_{23}$$

$$U(1)_Y$$

$$Z$$

$$B_{s,d} \rightarrow \mu\mu$$

$$b \rightarrow s + ll$$

$$\Delta B = 2$$

$$\text{EWPT}$$

$$pp \rightarrow ll$$

$$1/4 < b < 4, \quad 1/2 < t_\alpha < 2$$

$$m_{Z_{23}} \gtrsim 4 \div 5 \text{ TeV}$$

Near MFV in the  $b, t, \tau$  - sector

$$\text{E.g. } \mathcal{L} \approx \frac{g'^2}{m_{Z_{23}}^2} (U_{L3b} U_{L3q}^*)^2 \mathcal{O}_1^{bq} \quad q = s, d$$

$$\frac{m_{Z'_{23}}}{m_{Z_{23}}} = \frac{g_3}{g'} \frac{(1+b)t_\alpha}{\sqrt{b}(1+t_\alpha^2)}$$

$$b = \frac{|\phi|^2}{|\chi^l|^2 + |\chi^q|^2/9}$$

(an overall fit still lacking)

# Phenomenology at $\Lambda_{[12]}$ (universal)

$$U(1)_Y^{[3]} \times U(1)_{(B-L)/2}^{[12]} \times U(1)_{T_{3R}}^{[2]} \times U(1)_{T_{3R}}^{[1]}$$

$g_3 \qquad g_B \qquad g_T \qquad g_T$

$$g_3 \gg g_B, g_T$$

$$g_B = g' / \sin\alpha$$

$$g_T = g' / \cos\alpha$$

At  $\Lambda_{[12]}$  3 pars :  $m_{Z_{12}}$  and  $(\theta_R^{u,d})_{12}$

$$\mathcal{L} \approx \frac{g'^2}{m_{Z_{12}}^2} [(U_{R2s} U_{R2d}^*)^2 \tilde{\mathcal{O}}_1^{sd} + (U_{R2c} U_{R2u}^*)^2 \tilde{\mathcal{O}}_1^{cu}]$$

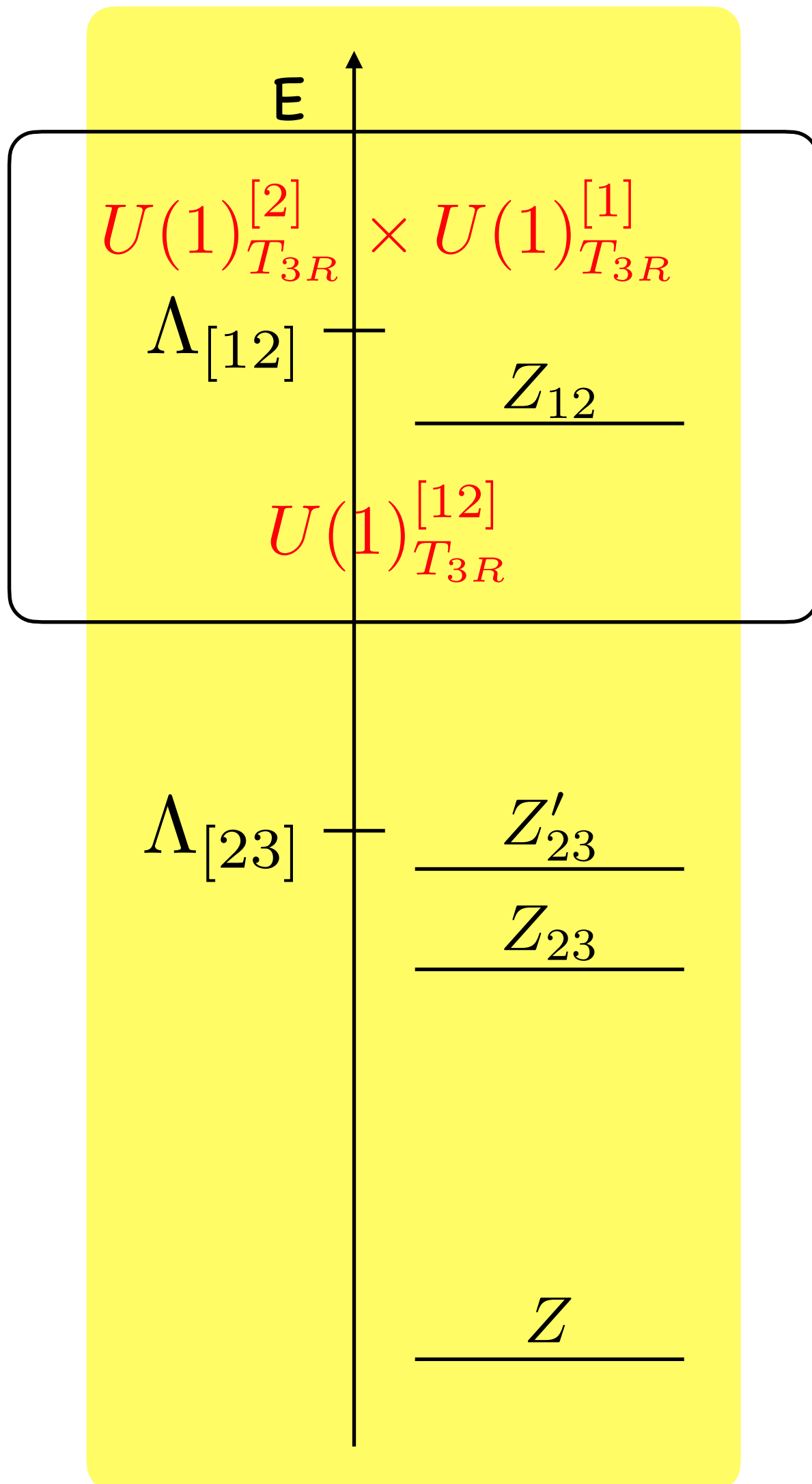
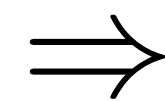
$$m_{Z_{12}} \gtrsim 100 \text{ TeV}$$

$$\Delta S = 2$$

$$\Delta C = 2$$

$$K \rightarrow \pi \nu \nu$$

$$\mu \rightarrow 3e$$





# Phenomenology (details)

$$g_3 \gg g_B, g_T \quad g_B = g' / \sin\alpha \quad g_T = g' / \cos\alpha$$

$\Psi$  = Standard fermions in interaction basis

$$\mathcal{L}_{Z_{23}} = g' Z_{23\mu} \bar{\Psi} \gamma_\mu \hat{Y}^{[23]} \Psi$$

$$\mathcal{L}_{Z'_{23}} = g_3 Z'_{23\mu} \bar{\Psi} \gamma_\mu Y^{[3]} \Psi$$

$$\hat{Y}^{[23]} = a_{12} Y^{[3]} + \cot\alpha \left( \frac{B-L}{2} \right)^{[12]} + t_\alpha T_{3R}^{[12]}$$

$$a_3 = \frac{b}{s_\alpha^2 c_\alpha^2 (1+b)^2}$$

$$a_{12} = \frac{\cot\alpha - bt_\alpha}{1+b}$$

$$b = \frac{|\phi|^2}{|\chi^l|^2 + |\chi^q|^2/9}$$

$$\mathcal{L}_{Z_{12}} = \frac{g'}{\cos\alpha} Z_{12\mu} \bar{\Psi} \gamma_\mu (T_{3R}^{[2]} - T_{3R}^{[1]}) \Psi$$

$$m_Z^2 = m_{Z_0}^2 + \delta m_Z^2$$

$$\delta m_Z^2 = -s_W^2 \frac{m_Z^4}{m_{Z_{23}}^2} (a_{12}^2 + a_3)$$

$$\mathcal{L}_Z = \mathcal{L}_Z^{SM} + g' s_W \frac{m_Z^2}{m_{Z_{23}}^2} Z_\mu \bar{\Psi} \gamma_\mu (a_{12} \hat{Y}^{[23]} + a_3 Y^{[3]}) \Psi$$

Back to the “mid-term” prospects  
of flavour

# Flavour precision tests, a partial list (2022)

Input	Reference	Measurement	UTfit Prediction	Pull
$\sin 2\beta$	[22], <b>UTfit</b>	0.688(20)	0.736(28)	-1.4
$\gamma$	[22]	66.1(3.5)	64.9(1.4)	+0.29
$\alpha$	<b>UTfit</b>	94.9(4.7)	92.2(1.6)	+0.6
$\varepsilon \cdot 10^3$	[38]	2.228(1)	2.00(15)	+1.56
$ V_{ud} $	<b>UTfit</b>	0.97433(19)	0.9738(11)	+0.03
$ V_{ub}  \cdot 10^3 \bullet$	<b>UTfit</b>	3.77(24)	3.70(11)	+0.25
$ V_{ub}  \cdot 10^3$ (excl)	[39]	3.74(17)		
$ V_{ub}  \cdot 10^3$ (incl)	[22]	4.32(29)		
$ V_{cb}  \cdot 10^3 \bullet$	<b>UTfit</b>	41.25(95)	42.22(51)	-0.59
$ V_{cb}  \cdot 10^3$ (excl)	<b>UTfit</b>	39.44(63)		
$ V_{cb}  \cdot 10^3$ (incl)	[40]	42.16(50)		
$ V_{ub} / V_{cb} $	[39]	0.0844(56)		
$\Delta M_d \times 10^{12} \text{ s}^{-1}$	[38]	0.5065(19)	0.519(23)	-0.49
$\Delta M_s \times 10^{12} \text{ s}^{-1}$	[38]	17.741(20)	17.94(69)	-0.30
$\text{BR}(B_s \rightarrow \mu\mu) \times 10^9$	[38]	3.41(29)	3.47(14)	-0.14
$\text{BR}(B \rightarrow \tau\nu) \times 10^4$	[38]	1.06(19)	0.869(47)	+0.96
$\text{Re}(\varepsilon'/\varepsilon) \times 10^4$	[38]	16.6(3.3)	15.2(4.7)	+0.27
$( q/p _D - 1) \times 10^2$		0.05(2.50)	0.8(4.0)	-0, 15
$\text{BR}(B^+ \rightarrow K^+ \nu\nu) 10^6$		23(7)	5.58(37)	+2.5
$\text{BR}(K^+ \rightarrow \pi^+ \nu\nu) 10^{11}$		10.6(4.0)	9.31(76)	+0, 3
$R_D$		0.344(26)	0.298(4)	+1.7
$R_{D^*}$		0.285(12)	0.254(5)	+2.3

UTfit Collaboration  
with some little integration

(No EDM's,  $\mu \rightarrow e\gamma$ , etc.)



# Current precision

Input	Reference	Measurement	UTfit Prediction	Pull
$\sin 2\beta$	[22], UTfit	0.688(20)	0.736(28)	-1.4
$\gamma$	[22]	66.1(3.5)	64.9(1.4)	+0.29
$\alpha$	UTfit	94.9(4.7)	92.2(1.6)	+0.6
$\varepsilon \cdot 10^3$	[38]	2.228(1)	2.00(15)	+1.56
$ V_{ud} $	UTfit	0.97433(19)	0.9738(11)	+0.03
$ V_{ub}  \cdot 10^3 \bullet$	UTfit	3.77(24)	3.70(11)	+0.25
$ V_{ub}  \cdot 10^3$ (excl)	[39]	3.74(17)		
$ V_{ub}  \cdot 10^3$ (incl)	[22]	4.32(29)		
$ V_{cb}  \cdot 10^3 \bullet$	UTfit	41.25(95)	42.22(51)	-0.59
$ V_{cb}  \cdot 10^3$ (excl)	UTfit	39.44(63)	Ciao.	
$ V_{cb}  \cdot 10^3$ (incl)	[40]	42.16(50)		
$ V_{ub} / V_{cb} $	[39]	0.0844(56)		
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current

4%th/exp

5%exp

5%exp

8%th

0.1%th

8%exp/th

2%exp/th

5%th/exp

4%th

9%exp

20%exp

30%th

100%th\*

35%exp

40%exp

8%exp

4%exp

with strong correlations!

# Conceivable progress in the “mid-term” of flavour

Input	Reference	Measurement	UTfit Prediction	Pull
$\sin 2\beta$	[22], UTfit	0.688(20)	0.736(28)	-1.4
$\gamma$	[22]	66.1(3.5)	64.9(1.4)	+0.29
$\alpha$	UTfit	94.9(4.7)	92.2(1.6)	+0.6
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$ V_{ub}  \cdot 10^3$ (excl)	[39]	3.74(17)		
$ V_{ub}  \cdot 10^3$ (incl)	[22]	4.32(29)		
$ V_{cb}  \cdot 10^3 \bullet$	UTfit	41.25(95)	42.22(51)	-0.59
$ V_{cb}  \cdot 10^3$ (excl)	UTfit	39.44(63)		
$ V_{cb}  \cdot 10^3$ (incl)	[40]	42.16(50)		
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$R_D$		0.344(26)	0.298(4)	+1.7
$R_{D^*}$		0.285(12)	0.254(5)	+2.3

current

4%th/exp

5%exp

5%exp

8%th

0.1%th

8%exp/th

2%exp/th

5%th/exp

4%th

9%exp

20%exp

30%th

100%th\*

35%exp

40%exp

8%exp

4%exp

mid-term

0.6%

0.8%

0.4%

1%

0.5%

2%

1.5%

4%

4%

30%\*

10%

20%

4%

2.5%

with strong correlations!

(No EDM's,  $\mu \rightarrow e\gamma$ , etc.)

# Summary

1. The flavour puzzle and the hierarchy problem are key open issues in BSM  
(and a strong motivation for the next HE collider)
2. Precision offers an indirect discovery potential of NP at MultiTeV, if any, before the next HE collider
3. If the Higgs is composite,  $\Lambda^f$  must be “low” and subject to flavour symmetries
4. “Deconstructing” SM gauge interactions may be a way to address the flavour puzzle and make sense of NMFV in  $b, t, \tau$  -physics



...as promised



Engadina, 1997





Zurich, around 2000



# Last but not least



where?, 2000

Cheers to Zoltan  
and to all his family